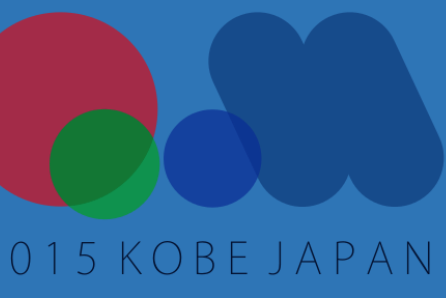


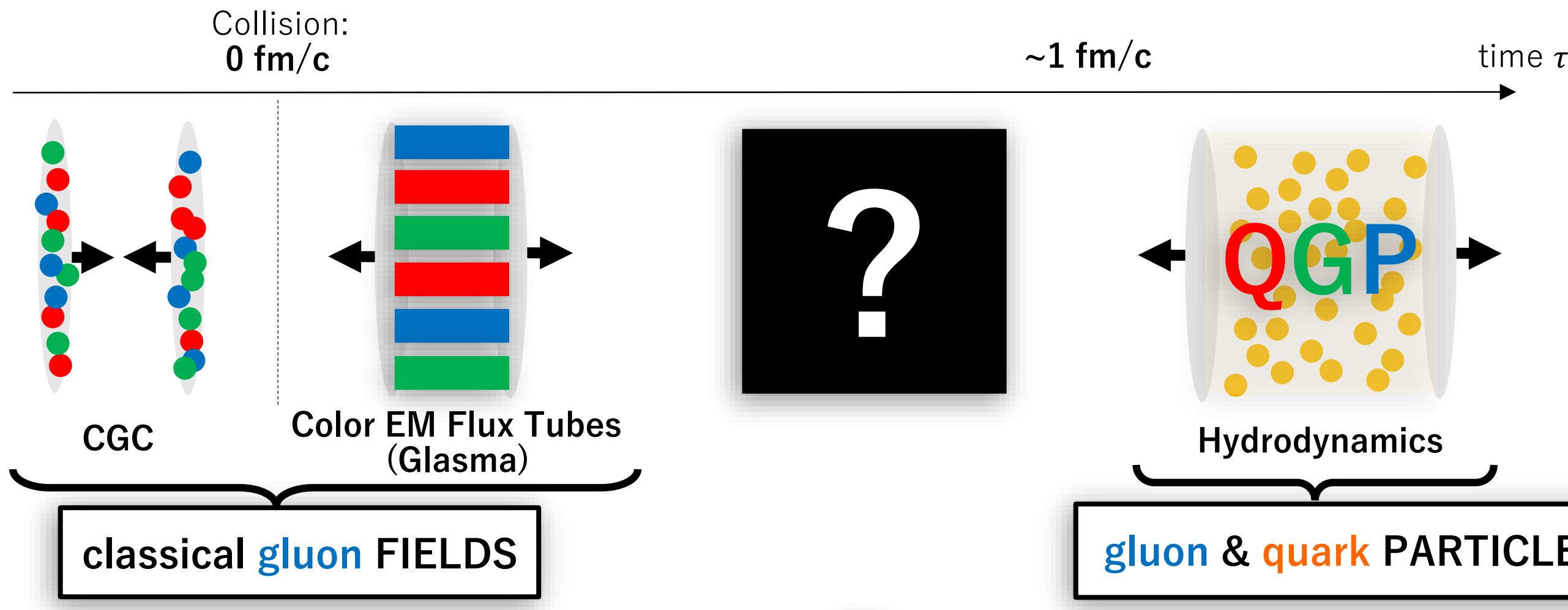
Quark Pair Production from Expanding Electromagnetic Flux Tube

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1. INTRODUCTION

1-1. Formation Process of QGP is a BIG QUESTION in HIC



Aim of this study | Try to reveal **?** focusing on the **quark** production based on the **Schwinger mechanism**

1-2. Previous Studies on Quark Production via Schwinger Mech.

Some of the important properties of the flux tubes are

1. **Longitudinal expansion** (Bjorken expansion) of the flux tubes
2. Presence of **magnetic fields** in parallel to electric fields,

whose effects on the Schwinger mech. are **NOT understood well** in previous studies:

	Expanding?	Magnetic Field?	Quark Mass	Backreaction?
Mihaila et al.	Yes	No	Limited: $\frac{m^2}{gE} = \frac{1}{4}$	Yes
Tanji	No	Yes (with arbitrary strength)	Arbitrary	Yes
Gelis et al.	Yes	Yes, but only for $\frac{gB}{gE} \sim 1$	Limited: $m \geq 60$ MeV	No

Our study **gets over** these limitations :

Our Study	Yes	Yes (with arbitrary strength)	Arbitrary	Yes
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2. THEORY

2-1. Formulation

Following **simplifications** are assumed hereafter:

1. **U(1) gauge theory with $N_f = 1$ fermion**
2. **Homogeneity in space: Boost invariance & transverse symmetry**
3. **Single flux tube with infinite length in the transverse direction**
4. **Longitudinal EM-field at initial time τ_0 : $\vec{E}(\tau_0) = (0, 0, E_0), \vec{B}(\tau_0) = (0, 0, B_0)$**

We solve the EoMs (non-perturbatively w.r.t the strong fields gA_μ):

$$0 = [i\gamma^\mu(\partial_\mu + igA_\mu) - m]\psi, \quad F^{\nu\mu}{}_{;\mu} = g\bar{\psi}\gamma^\nu\psi$$

and **canonically quantize $\psi \rightarrow \hat{\psi}$ at each instant of (proper) time τ** :

$$\hat{\psi}(x) = \sum_s \sum_n \int d\lambda \left[\psi_{n,\lambda,s}^{(+)}(x_T, \tau) \hat{a}_{n,\lambda,s}(\tau) + \psi_{n,\lambda,s}^{(-)}(x_T, \tau) \hat{b}_{n,-\lambda,s}^\dagger(\tau) \right] \frac{e^{i\lambda\eta}}{\sqrt{2\pi}}$$

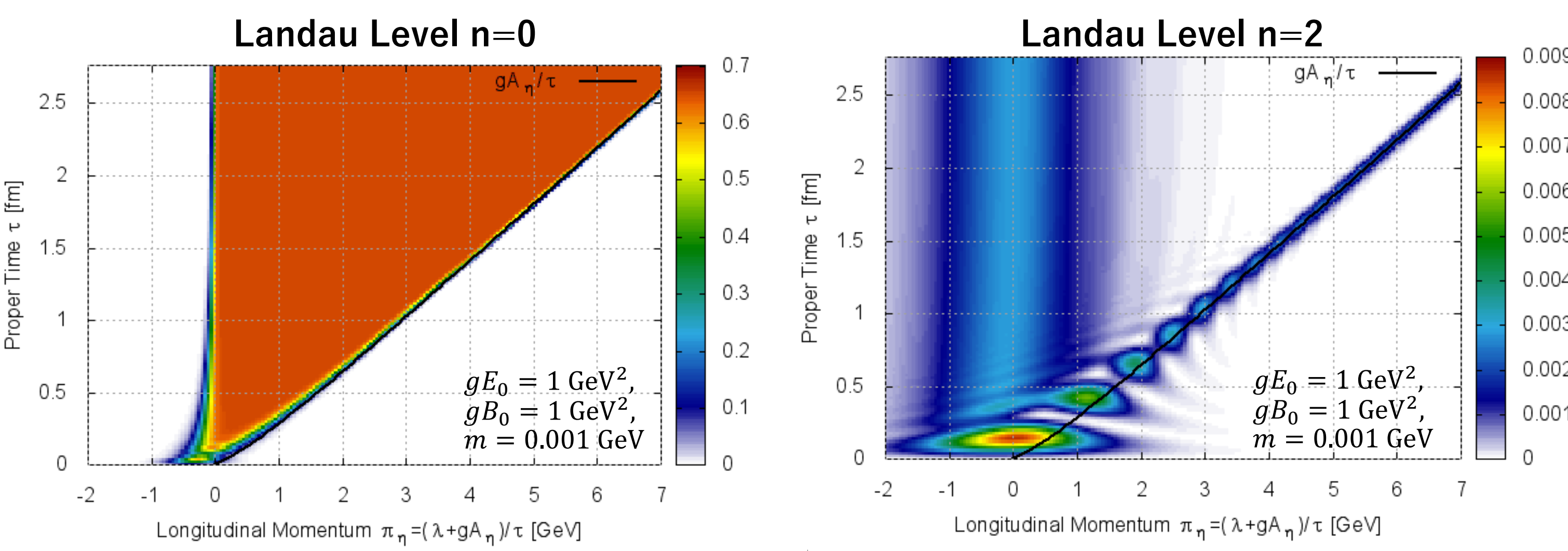
REMARK 1: 3 labels = **spin s** , **Landau level n** , and **(canonical) longitudinal momentum λ** .

REMARK 2: The instantaneous mode functions $\psi_{n,\lambda,s}^{(\pm)}(x_T, \tau)$ are defined by the approximate sol. of the Dirac eq. under an adiabatic assumption: $\partial_\tau A^\mu = 0$

Tanji; Schmidt, Blashke, Smolyansky ...; Kluger, Mottola, Eisenberg ...

3. RESULT I – WithOUT Backreaction $g=0$ (Analytical)

3-1. Longitudinal Distribution $d^2N/dy d\lambda/S_\perp$ [fm^{-2}]



The distribution is saturated because of Pauli's principle.

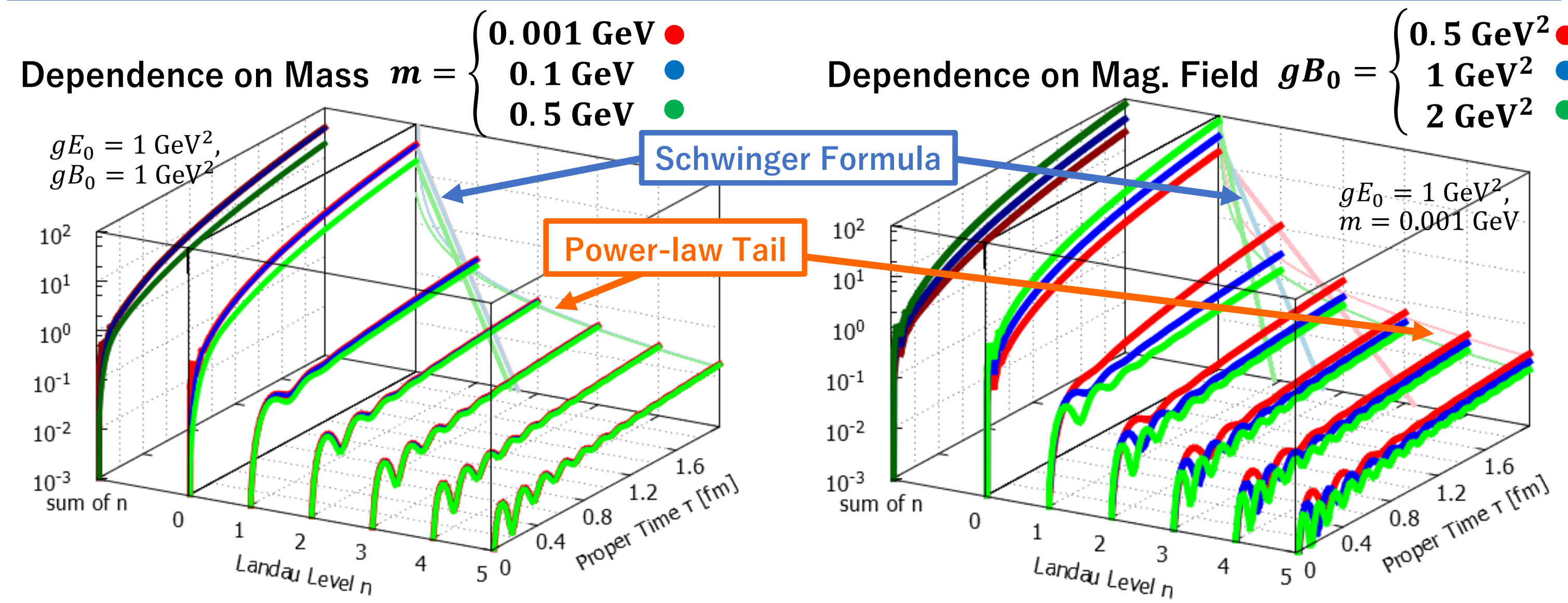
- Particles are created constantly with 0 longitudinal momentum.
- After being created, particles are accelerated by electric fields in the long. direction obeying the classical EoM: $0 = d\lambda/d\tau$

2-peak structure: **I** and **II**.

I: Particles which annihilate w/ their pair immediately after their creation. This peak depends on the choice of the instantaneous particle picture.

I has pwr-dep. on n : $\sim ((m^2 + 2ngB)/gE)^{-2}$

3-2. Number Density $dN/dy/S_\perp$ [fm^{-2}]



The density is consistent w/ **Schwinger's formula** $\propto \exp[-\pi(m^2 + 2ngB)/gE]$ for small n (m_T), while it has a **power-law tail** $\propto [(m^2 + 2ngB)/gE]^{-2}$ for large n (m_T). cf. Mihaila et al. (2009)

Lowest Landau contribution ($n=0$) dominates the particle production.

Larger particle production $\propto gB$ for increasing gB .

Very fast quark production: For example, $dN/dy/S_\perp \sim 1 \text{ fm}^{-2}$ particles are produced at $\tau \sim 0.2 \text{ fm}$, which roughly corresponds to $dN_q/dy \sim N_c \times N_f \times 1 \times S_\perp \sim 1000$ in HIC. cf. Gelis et al. (2006)

Combining with the experimental fact $dN_h/dy \lesssim 1000$, **classical gluon fields should have decayed very fast $\tau \lesssim 0.2 \text{ fm}$** , otherwise we would have too huge # of particles.

5. SUMMARY

What we did

The **quark** production in the early stage dynamics of HIC based on the **Schwinger mechanism** is extensively studied.

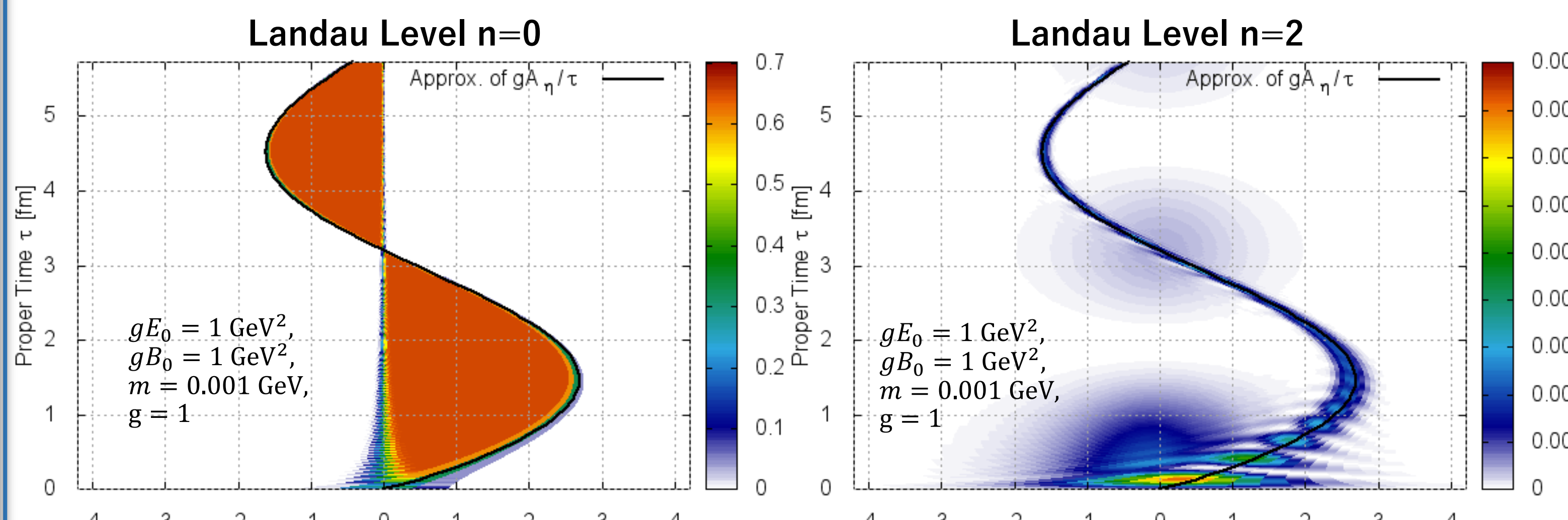
What we have learned

1. Effects of **the longitudinal expansion** and **the longitudinal magnetic fields** on the Schwinger mech., which are the missing pieces of previous studies, are clarified.
2. The Schwinger mech. can explain a **fast quark production** and, combining w/ the experimental result on the multiplicity, it suggests that **the classical gluon fields should have decayed very fast** (though we did not say anything about its physics).
3. The **backreaction effects from the produced quarks** are also investigated, and an approximate formula is obtained in the LLL approximation and the massless limit.

There are several topics (energy density, pressure, anomalous current, etc) which I could not cover in this poster. If you are interested in these topics, feel free to ask me!

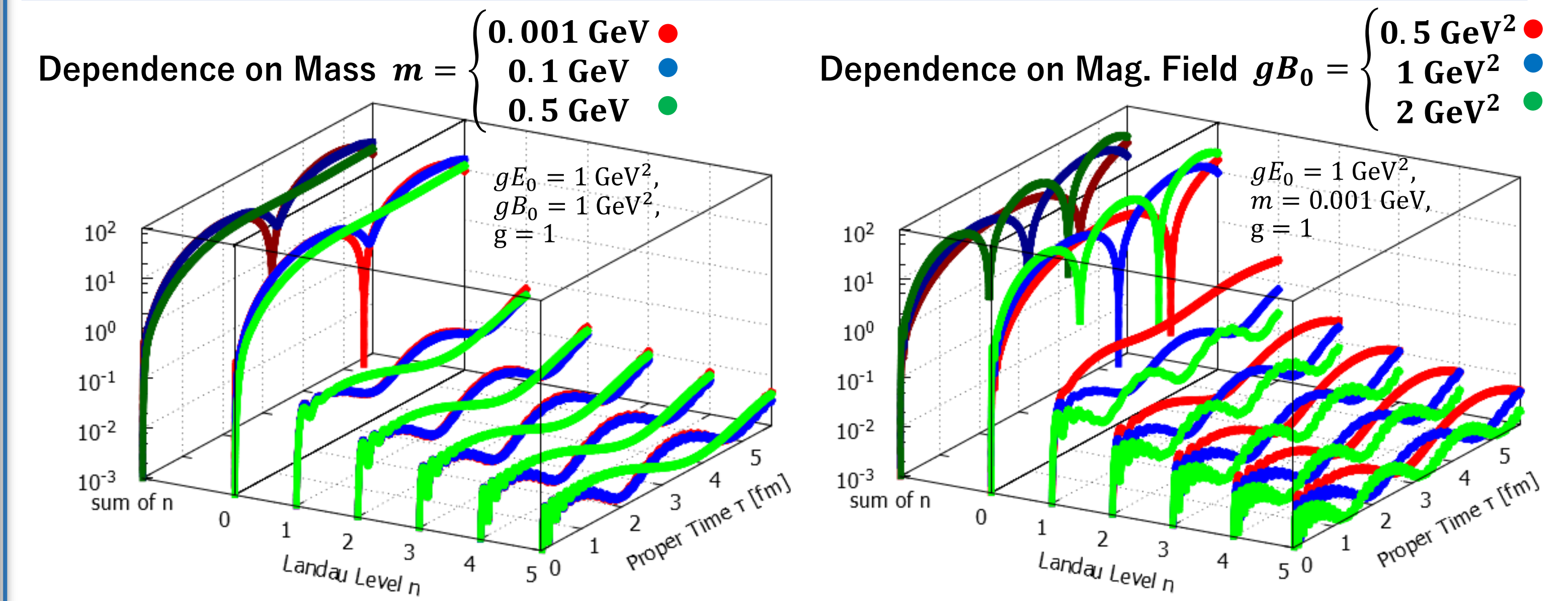
4. RESULT II – With Backreaction $g \neq 0$ (Numerical)

4-1. Longitudinal Distribution $d^2N/dy d\lambda/S_\perp$ [fm^{-2}]



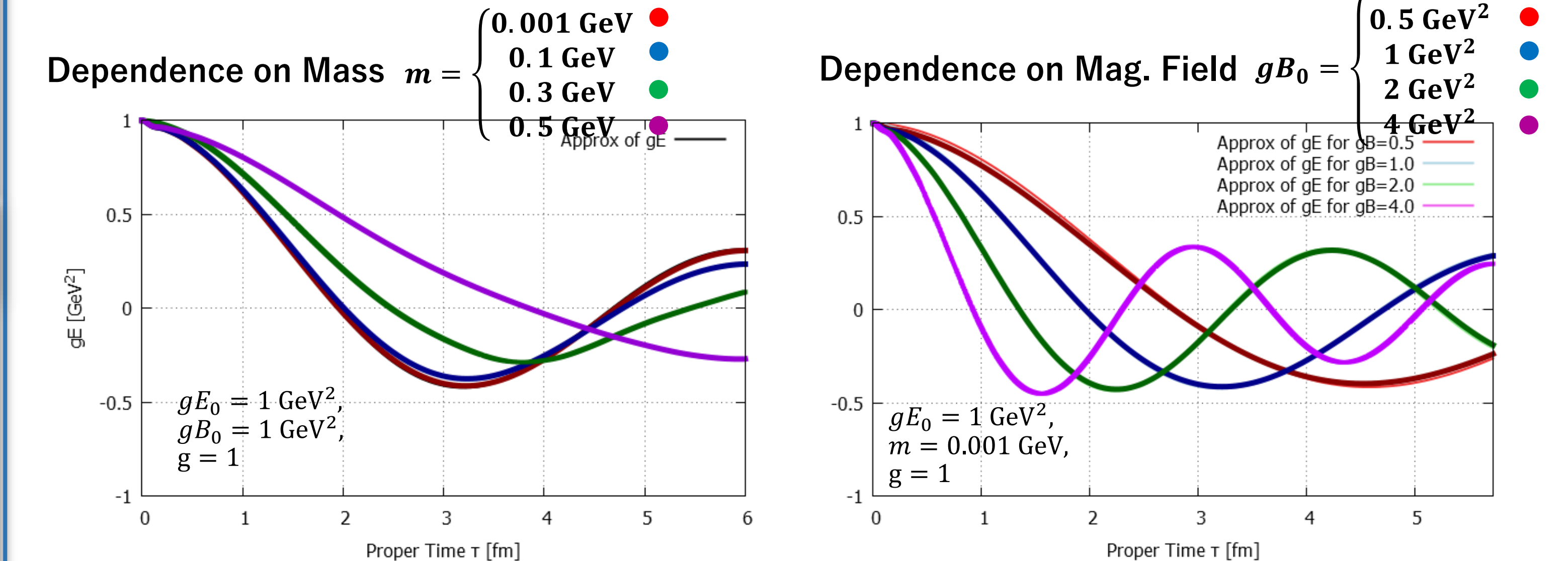
- Plasma oscillation** with timescale $\tau_{\text{osc}} \sim 1/\sqrt{\frac{g^2 N_f N_c}{2\pi^2} gB}$ occurs as a result of the backreaction.
- In addition to the oscillating behavior, the longitudinal momentum π_η decreases as $\sim 1/\sqrt{\tau}$ reflecting the longitudinal expansion of the system.

4-2. Number Density $dN/dy/S_\perp$ [fm^{-2}]



- The oscillation is "flattened" with increasing the mass.
- Faster oscillation with increasing gB .
- Effect of backreaction is small for small values of $\tau \lesssim 0.2 \text{ fm}$, while it is not for large values of $\tau \gtrsim 0.2 \text{ fm}$, where $dN/dy/S_\perp \propto \sqrt{\tau}$ (τ^2) with (without) backreaction.

4-3. Decay of Electric Fields



- gB never decays:** Rotational sym. always holds due to the simplification 2&3 (i.e., $j^\theta = 0$)
- Faster decay and slower oscillation with increasing mass.
- Faster oscillation for stronger gB .
- Long. expansion makes the field decay faster $\propto 1/\sqrt{\tau}$ compared to the non-expanding case.

Note: Approximate Solution for the Backreaction Problem

If we take LLL approx. and massless limit, we can **analytically** solve the backreaction eqs. to find

$$E(\tau) = E_0 \left[\frac{J_1(s\tau_0)}{J_1(s\tau_0)Y_0(s\tau_0)} Y_0(s\tau) - \frac{Y_1(s\tau_0)}{Y_1(s\tau_0)J_0(s\tau_0)} J_0(s\tau) \right], \quad \frac{1}{s_\perp} \frac{dN}{dy}(\tau) = \frac{gE_0}{2} s\tau \left[\frac{J_1(s\tau_0)}{J_1(s\tau_0)Y_0(s\tau_0)} Y_1(s\tau) - \frac{Y_1(s\tau_0)}{Y_1(s\tau_0)J_0(s\tau_0)} J_1(s\tau) \right]$$

where $s^2 = \frac{g^2}{2\pi^2} gB$ (which roughly corresponds to $s^2 \sim \frac{N_c N_f g^2}{2\pi^2} gB$ in QCD).