

Introduction to strong-field QED

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(an SPDR in iTHEMS, studying high-energy particle & nuclear physics)

Plan

(QED = quantum electrodynamics)

Brief introduction to **Strong-field QED**

= an area to study **what happens by a super strong light**

(= electromagnetic field)

PART I: Basics of strong-field QED

- Why interesting \Rightarrow Unexplored non-perturbative regime of physics
Nontrivial phenomena that we've never seen e.g. Vacuum decay, Birefringence, New phase of matter, ...
- Why timely \Rightarrow Recent availability of strong fields e.g. High-power laser, Magnetar, heavy-ion collisions, ...

PART II: Sauter-Schwinger effect and its connection to other areas

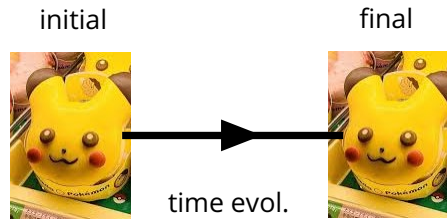
- An interdisciplinary topic
 \Rightarrow Exchange of ideas/techniques among physics is quite useful for better understandings
 \Leftarrow An example from my experience: High-harmonic generation from the vacuum
 \Leftarrow Math. & cond-mat. ideas helped me a lot

[[HT](#), Fujimori, Misumi, Nitta, Sakai, (2021)] [[HT](#), Hongo, Ikeda, (2021)]

PART I: Basics of strong-field physics

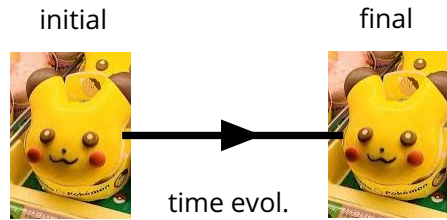
PART II: Sauter-Schwinger effect and its connection to other areas

What if field becomes strong ?



No field

What if field becomes strong ?



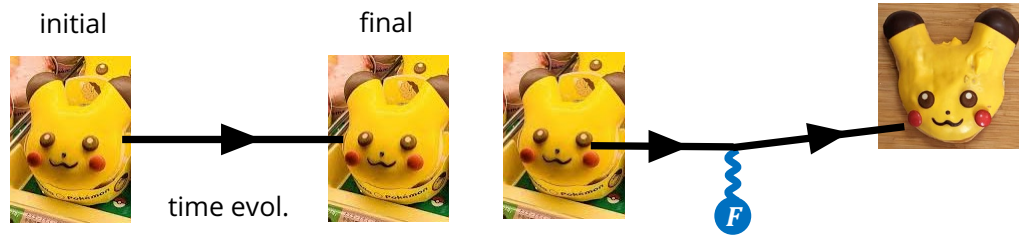
No field

Weak field

Strong field



What if field becomes strong ?



No field

Weak field

Strong field

Only minor changes

⇒ Perturbative

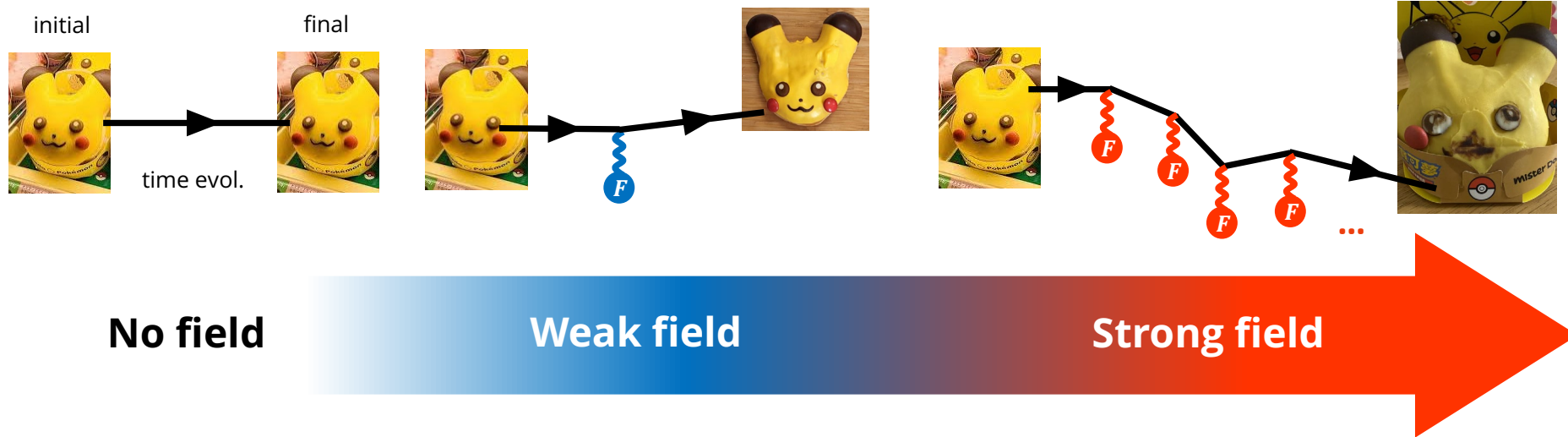
⇒ Very well understood
in both exp.& theor.

ex.) Electron (anomalous) magnetic moment $a := \frac{g-2}{2}$
≈ Electron energy shift in a weak magnetic field

$$a(\text{theor.}) = 1159652182.03 \dots \times 10^{-12}$$

$$a(\text{exp.}) = 1159652180.73 \dots \times 10^{-12} \quad [\text{Aoyama, Kinoshita, Nio (2017)}]$$

What if field becomes strong ?



Only minor changes

⇒ Perturbative

⇒ Very well understood
in both exp.& theor.

Big change !

⇒ Non-Perturbative

⇒ Not understood well

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If field becomes strong, physics becomes totally different & nontrivial

When is field “strong” ?

Strong-field condition:

To significantly modify the original system with typical energy Δ , the field must be more energetic than Δ

\Rightarrow Strong-field condition: $\Delta < (\text{energy scale of the field})$

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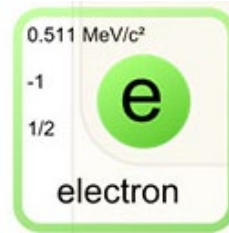
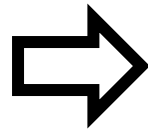
⇒ **Strong-field condition: $\Delta < (\text{energy scale of the field})$**

Estimate of the minimum field strength:

Standard Model

The (matter) particle having the minimum energy

mass = 12.3 MeV/c ² charge = 2/3 spin = 1/2	mass = 1.275 GeV/c ² charge = 2/3 spin = 1/2	mass = 173.2 GeV/c ² charge = 2/3 spin = 1/2	mass = 120 MeV/c ² charge = 1 spin = 1	mass = 125 GeV/c ² charge = 0 spin = 0
u up	c charm	t top	g gluon	H Higgs boson
mass = 4.2 MeV/c ² charge = -1/3 spin = 1/2	mass = 95 MeV/c ² charge = -1/3 spin = 1/2	mass = 4.18 GeV/c ² charge = -1/3 spin = 1/2	mass = 0 charge = 0 spin = 1	mass = 91.1876 GeV/c ² charge = 0 spin = 0
d down	s strange	b bottom	γ photon	Z Z boson
mass = 0.511 MeV/c ² charge = -1 spin = 1/2	mass = 105.658 MeV/c ² charge = -1 spin = 1/2	mass = 1.777 GeV/c ² charge = -1 spin = 1/2	mass = 91.1876 GeV/c ² charge = 0 spin = 1	mass = 80.379 GeV/c ² charge = 0 spin = 1
e electron	μ muon	τ tau	Z Z boson	W W boson
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ν _e electron neutrino	ν _μ muon neutrino	ν _τ tau neutrino	W W boson	W W boson



- Charged ⇒ Couples to electromagnetic (EM) field
- $eE \text{ MeV}^2 > m_e^2 = (0.511 \text{ MeV})^2 \approx 0(10^{28} \text{ W/cm}^2)$

∴ **EXTREMELY strong ⇒ Impossible to realize within the current technology**



LED $\approx 10^{-5} \text{ W/cm}^2$



Glasses $\approx 10^3 \text{ W/cm}^2$



Laser beam machine $< 10^{10} \text{ W/cm}^2$

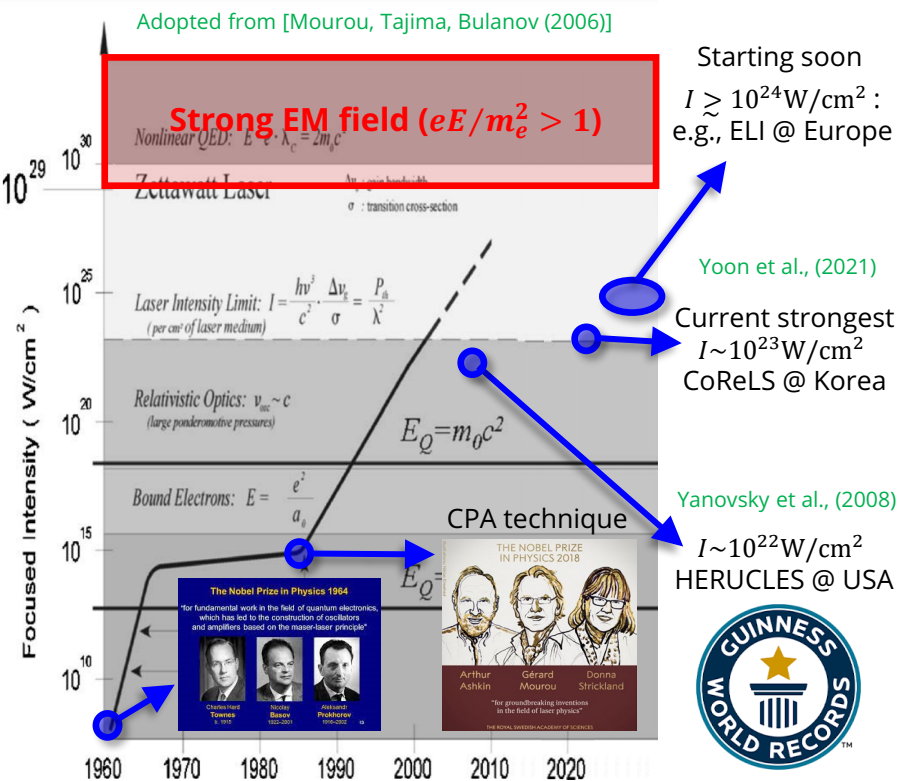
Recent availability of strong EM fields

The situation is changing:

Becoming able to create/observe strong EM fields

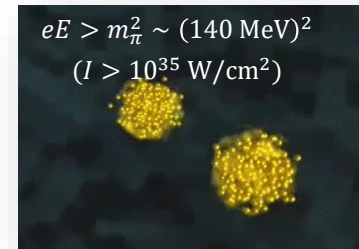
⇒ NOW is the best time to study physics of strong EM field
 (=: strong-field QED)

High-power laser



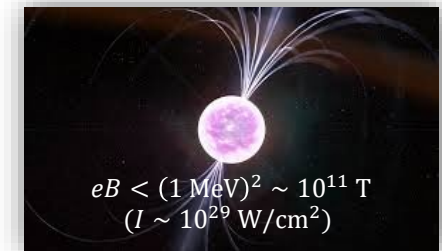
Systems with extreme conditions

- Heavy-ion collisions



At RHIC (2000~), LHC (2012~)

- Magnetar



Observation by IXPE, XL-Calibur (2022~)

Press release by Tamagawa lab. (2022)

What can happen with strong EM field ?

When $eE > m_e^2$, many non-trivial phenomena have been predicted to occur:

Review: [Fedotov, Ilderton, Karbstein, King, Seipt, HT, Torgrimsson, Phys. Rept. (2023)]

Sauter (1931), Schwinger (1951)

What I like the most: Production of particles from the vacuum (Sauter-Schwinger effect)

A few examples among many others (don't explain; ask later if interested)

Birefringence of photon in vacuum

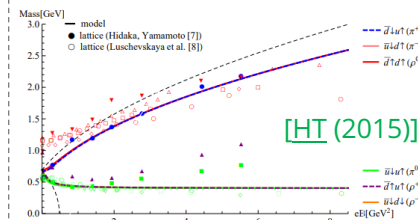
Photon splitting



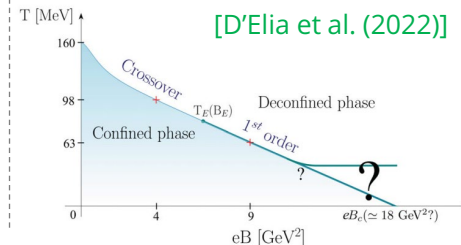
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Hadron mass change (for $eB > \Lambda_{\text{QCD}}^2$)



QCD phase diagram (for $eB > \Lambda_{\text{QCD}}^2$)



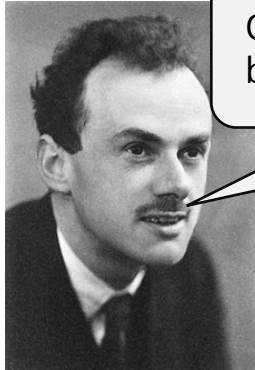
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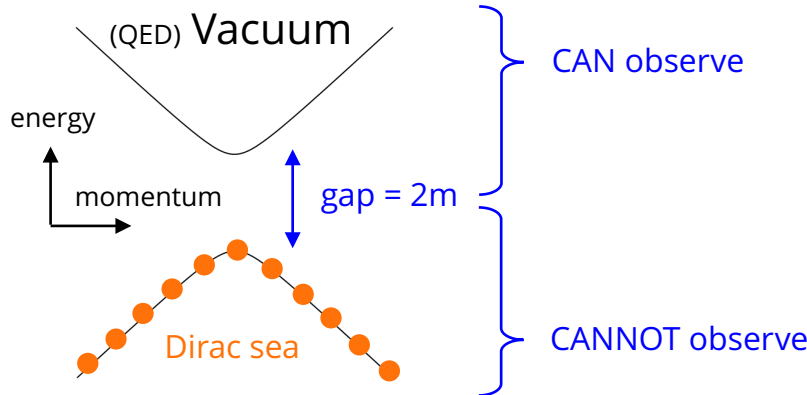
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Dirac (1930)

Our vacuum is **not an empty space**, but has a structure **similar to semi-conductor**



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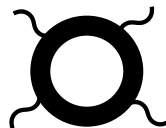
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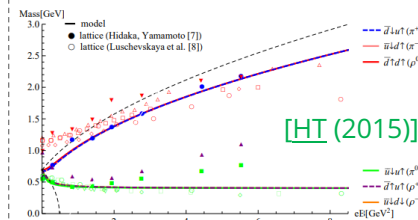
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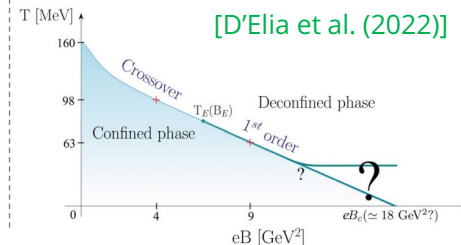
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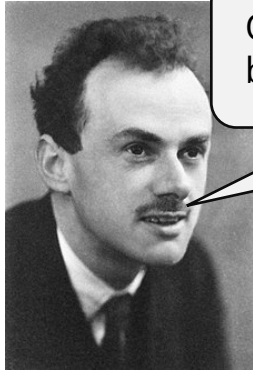
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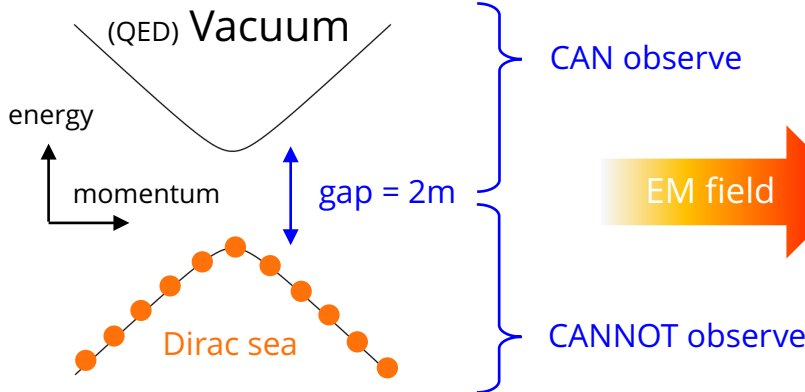
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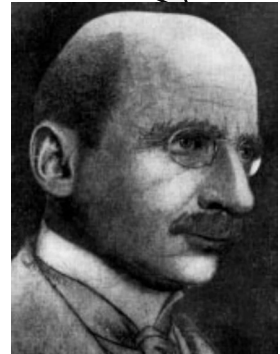
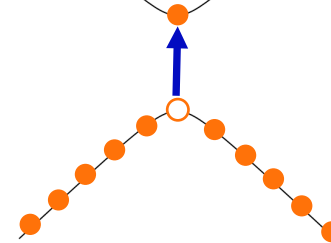
Dirac (1930)

Our vacuum is **not an empty space**, but has a structure **similar to semi-conductor**



Then, **real particles will be produced out of the vacuum** by strong enough field

Particle production !



Sauter (1931)

A few examples among many others (don't explain; ask later if interested)

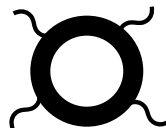
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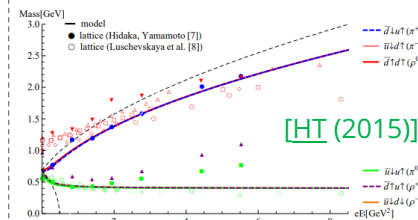
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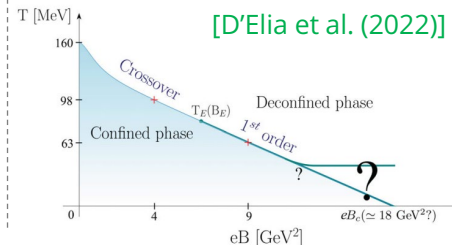
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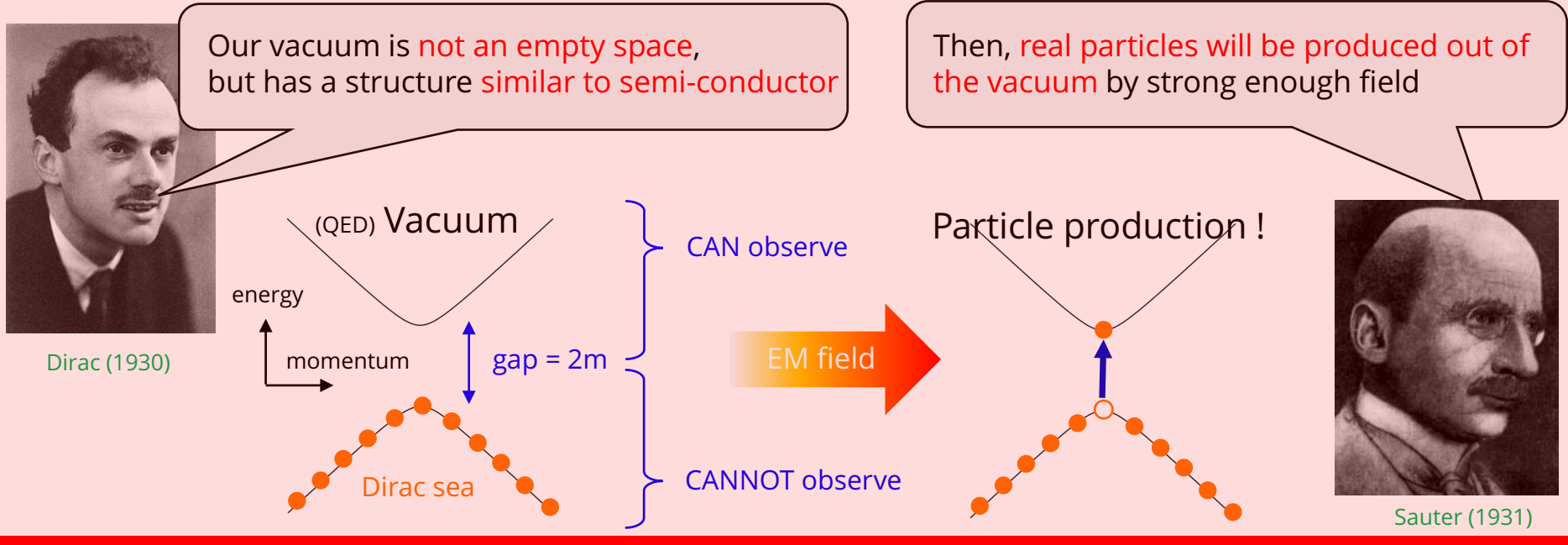
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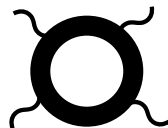
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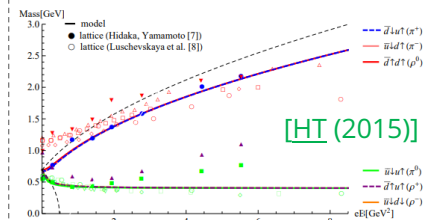
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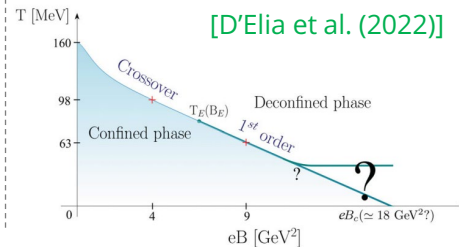
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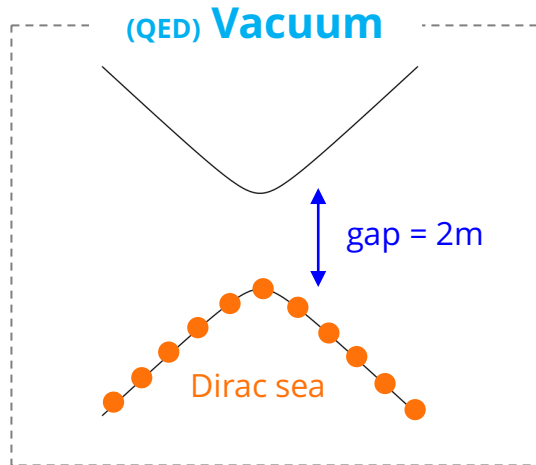
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**Similar particle production mechanisms
appear in many other areas of physics**

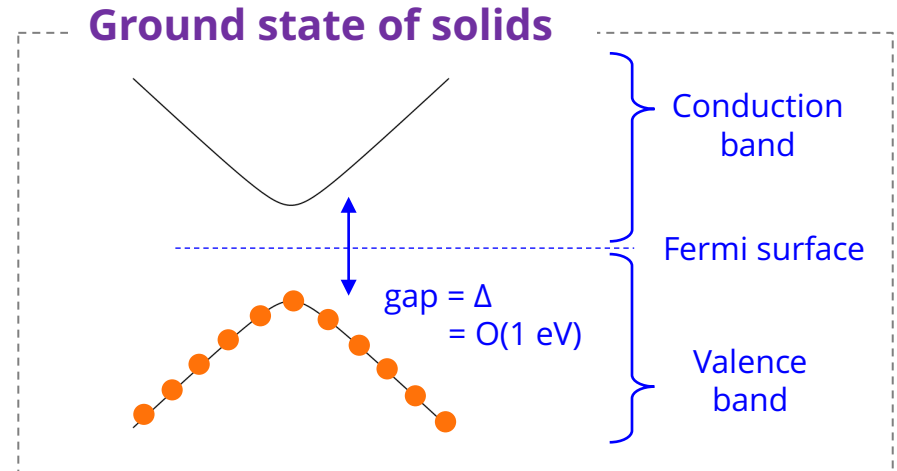
Similar particle production mechanisms appear in many other areas of physics

The most prominent example: Landau-Zener transition in solid

Landau (1932), Zener (1932)



v.s.



- The responses of the vacuum and ground state by an external EM field should be similar
- Indeed, the particle prediction rate and the excitation rate obey the same exponential formula

$$\Gamma_{\text{Sauter-Schwinger}} = \# \exp\left[\# \frac{m^2}{eE}\right] \quad \Gamma_{\text{Landau-Zener}} = \# \exp\left[\# \frac{\Delta^2}{eE}\right]$$

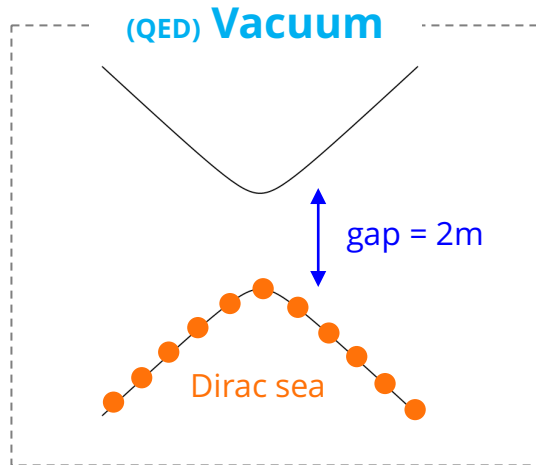
Note: Similar exponential formula holds for particle production by other strong fields as well

Strong gravitational field \Rightarrow Hawking radiation, Strong inflaton field \Rightarrow (p)reheating of the early Universe, ...

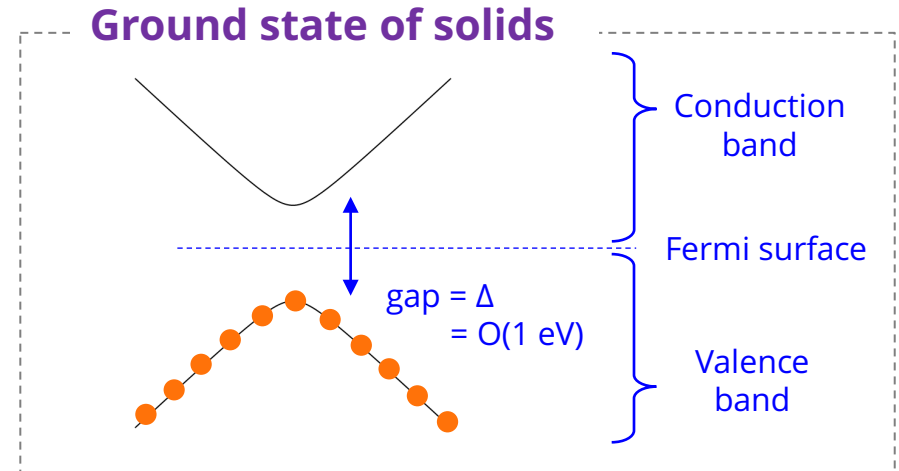
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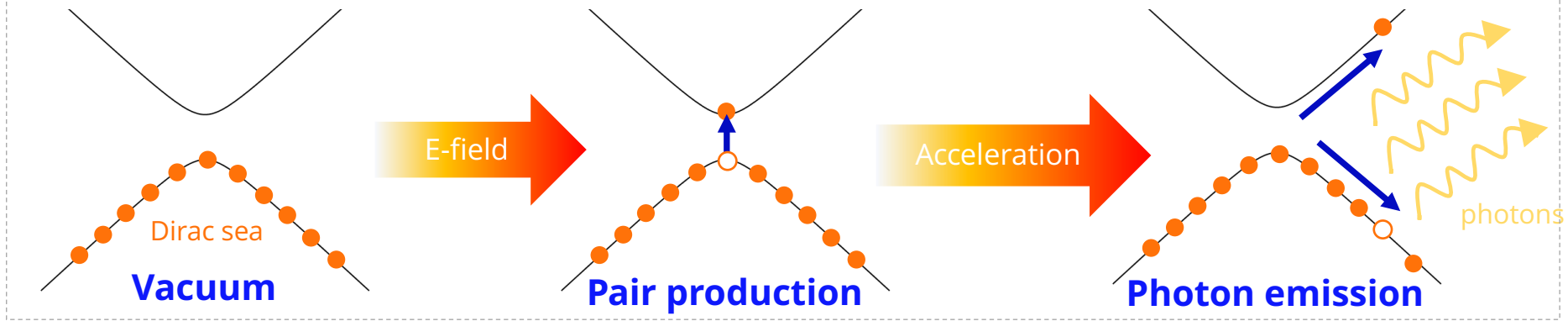
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- Sauter-Schwinger effect has not been verified yet, but Landau-Zener transition has been observed and utilized, e.g., to design devices.

\Rightarrow Importing (exporting) ideas from (to) other areas of physics is quite useful to better understand the Sauter-Schwinger effect (or strong-field QED in general)

Example: High-harmonic generation from the vacuum

Pair production by E-field leads to emission of photons

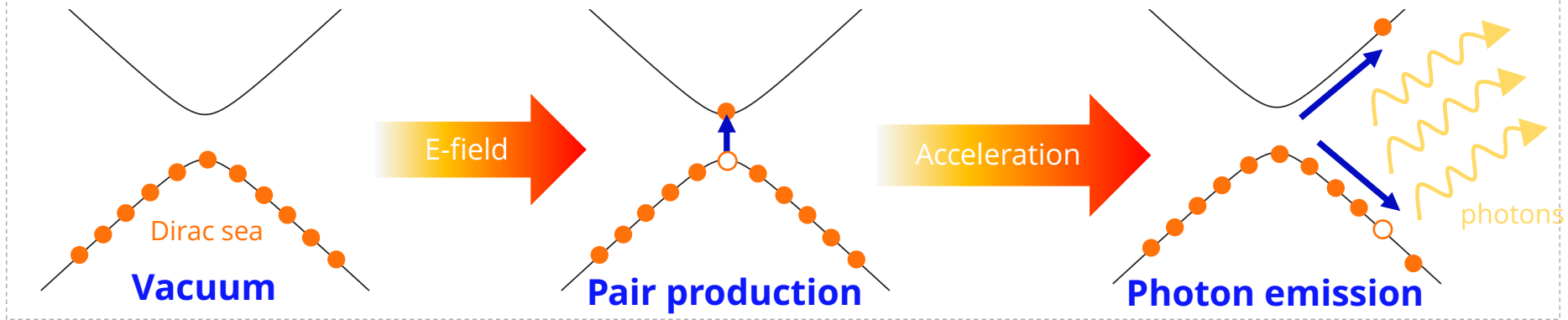


✓ Consider AC field with frequency Ω :

- Very naïve expectation: photon has frequency ω same to AC field $\omega = \Omega$

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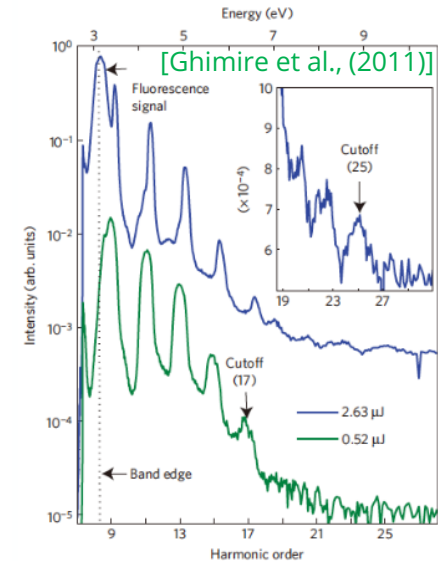
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✓ Recent observation in semi-conductors

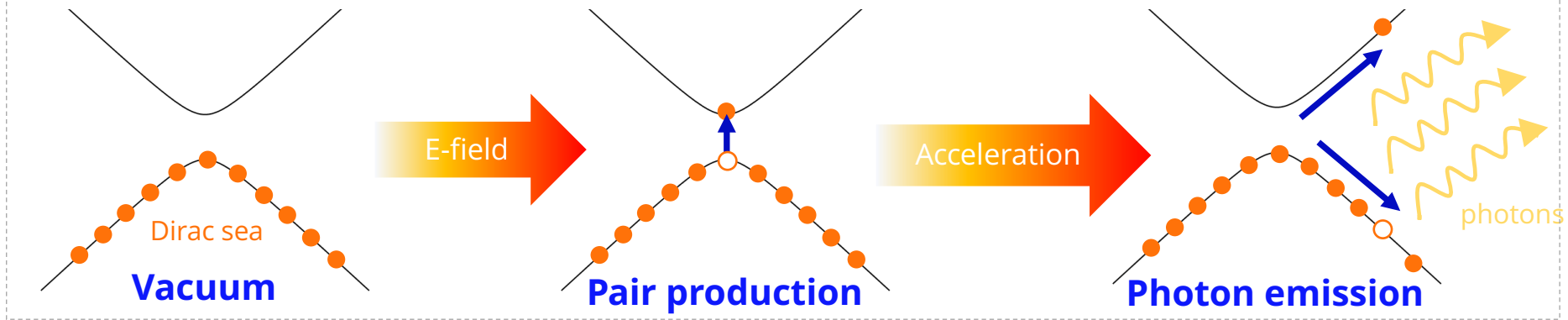
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- Naïve expectation is wrong
⇒ Exp. observation: Plateau structure
- Theoretical formulation is still immature even in cond.-mat
⇒ mostly numerics, and analytical understanding is lacking



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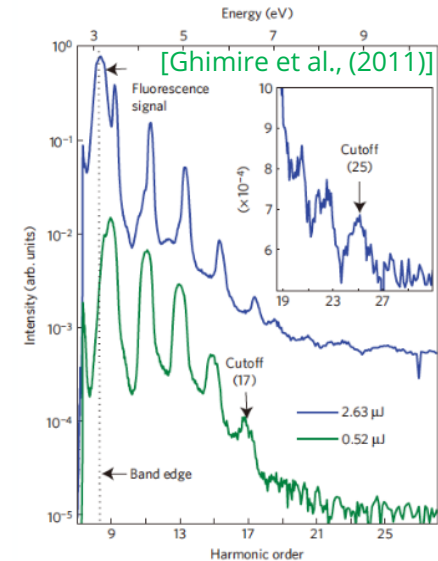
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Q1: High-harmonic generation in QED from the vacuum ?

Q2: Is there any nice analytical method ?

A: Yes. Mathematics helped me to answer those !

How to formulate ? Typical math of strong-field QED

Roughly speaking, strong-field-QED calculations is reduced to the following:

Step 1: Solve a differential equation (to obtain wavefunction $\Phi \approx$



The simplest example: Klein-Gordon equation with a time-dependent potential

$$0 = [\hbar^2 \partial_t^2 + Q(t)]\phi(t)$$

something small
(Planck constant)

Potential
(= info. of strong field)

wavefunction

decode

Step 2: Analyze the wavefunction (to predict sthg measurable $O[\Phi] \approx$



- O depends on what you want to measure
- Typically, a bi-linear $\phi^\dagger X \phi$

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Solve a differential equation is purely a mathematical problem
⇒ Any good idea from mathematics
⇒ Exact WKB

Exact WKB method

Exact WKB = a nice method to solve ODE with a small parameter

[Voros (1983)]

[Pham, Dillinger, Delabaere,

Aoki, Koike, Takei, ...]

= **"usual" WKB** + **Borel resum.**

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A perturbation theory w.r.t. \hbar (or adiabatic approx.)

• Consider $0 = [\hbar^2 \partial_t^2 + Q(t)]\phi(t) \stackrel{t \equiv \hbar\tau}{\Leftrightarrow} [\partial_\tau^2 + Q(\hbar\tau)]\phi(\tau)$

$$\Rightarrow \phi_{\pm}(t; \hbar) := \exp\left[\mp \frac{i}{\hbar} \int_{t_0}^t dt' \sqrt{Q(t')}\right] \times \sum_{n=0}^{\infty} \psi_{\pm,n}(t) \hbar^n$$

0th order = plane wave

$$\sim \exp\left[\mp \frac{i}{\hbar} \sqrt{Q} t\right]$$

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Perturbation w.r.t. \hbar

- WKB expansion makes sense if the **perturbative part** is convergent
- However, $\psi_{\pm,n} \sim n!$ in general (e.g., Airy function $Q(t) \propto t$)

\Rightarrow WKB expansion has zero radius of convergence \Rightarrow ill-defined !!!

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A resummation scheme for factorially divergent $\sim n!$ series

• Consider the div. part of WKB expansion $\psi_{\pm}(t; \hbar) := \sum_{n=0}^{\infty} \psi_{\pm,n}(t) \hbar^n$

① Construct "Borel transformation": $B[\psi_{\pm}](t; \eta) := \sum_{n=0}^{\infty} \frac{\psi_{\pm,n}(t)}{n!} \eta^n$

② Laplace trans. gives "Borel sum": $\Psi_{\pm}(t; \hbar) := \int_0^{\infty} \frac{d\eta}{\hbar} e^{-\eta/\hbar} B[\psi_{\pm}](t; \eta)$

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② Laplace trans. gives "Borel sum": $\Psi_{\pm}(t; \hbar) := \int_0^{\infty} \frac{d\eta}{\hbar} e^{-\eta/\hbar} B[\psi_{\pm}](t; \eta)$

• Ψ_{\pm} is well-defined and is a natural analytic continuation of ψ_{\pm}

⇐ Indeed, reduces to the original result if you admit $\sum \int = \int \sum$

$$\Psi_{\pm} = \int_0^{\infty} \frac{d\eta}{\hbar} e^{-\eta/\hbar} \sum_{n=0}^{\infty} \frac{\psi_{\pm,n}(t)}{n!} \eta^n \sim \sum_{n=0}^{\infty} \frac{\psi_{\pm,n}(t)}{n!} \int_0^{\infty} \frac{d\eta}{\hbar} e^{-\eta/\hbar} \eta^n = \sum_{n=0}^{\infty} \psi_{\pm,n}(t) \hbar^n$$

⇒ **Ψ_{\pm} gives a well-defined version of the WKB solution !**

• Note: in practice, some approximations shall be used in Borel resum...

How to formulate ? Typical math of strong-field QED

Step 1: Solve a differential equation (to obtain wavefunction $\Phi \approx$



The simplest example: Klein-Gordon equation with time-dependent potential

$$0 = [\hbar^2 \partial_t^2 + Q(t)]\phi(t)$$

something small
(Planck constant)

potential

wavefunction



decode

Step 2: Analyze the wavefunction (to predict sthg measurable $O[\Phi] \approx$



- O depends on what you want to measure
- Typically, a bi-linear $\phi^\dagger X \phi \Leftarrow$ For HHG, current $J \sim e\phi^\dagger \hat{p} \phi$ or $J \sim \bar{\psi} \gamma \psi$ (for fermions)

Application to high-harmonic generation (1/2)

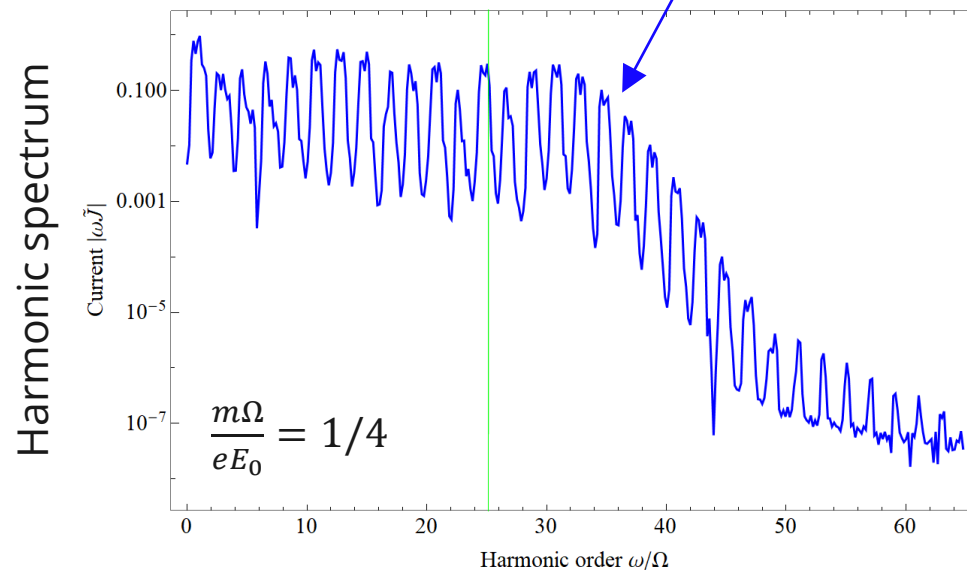
Setup: A monochromatic E-field $E(t) = E_0 \cos(\Omega t)$ applied onto QED vacuum

What I did: Compute harmonic spectrum **numerically** and **analytically with WKB**

Application to high-harmonic generation (1/2)

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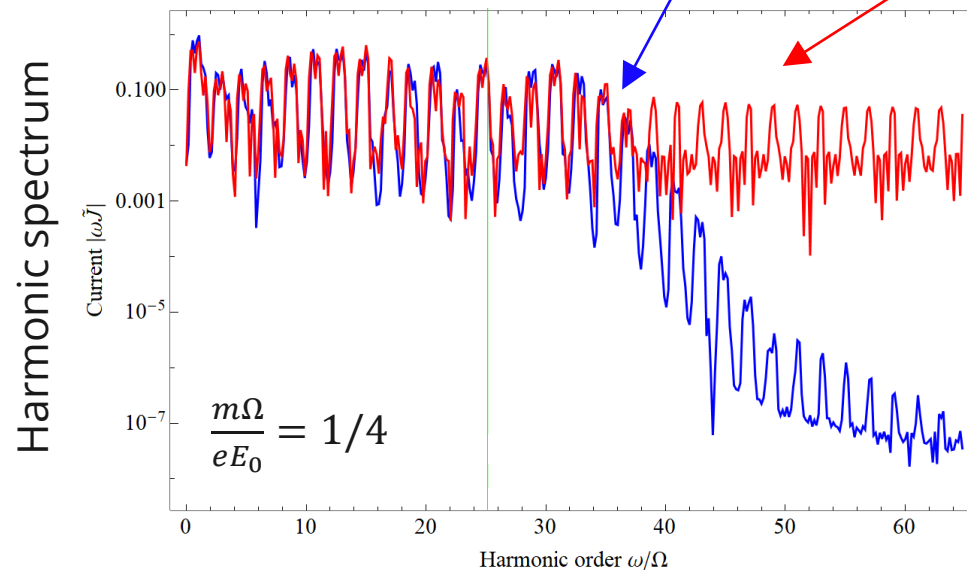


- High harmonics appears also in QED when field becomes strong !

Application to high-harmonic generation (1/2)

Setup: A monochromatic E-field $E(t) = E_0 \cos(\Omega t)$ applied onto QED vacuum

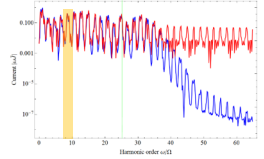
What I did: Compute harmonic spectrum **numerically** and **analytically with WKB**



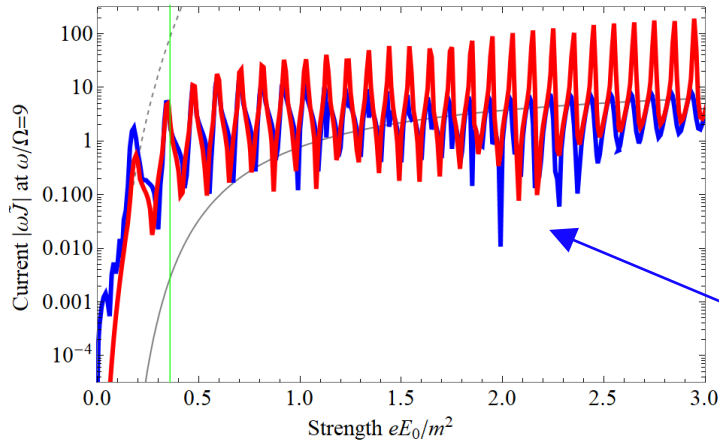
- High harmonics appears also in QED when field becomes strong !
- WKB is good in the plateau regime (i.e., before the spectrum drops)

Application to high-harmonic generation (2/2)

- WKB works more in the deep non-perturbative regime $E_0 \rightarrow$ large, $\Omega \rightarrow$ small
- Demonstration: magnitude of the harmonic peak at $\omega/\Omega = 9$

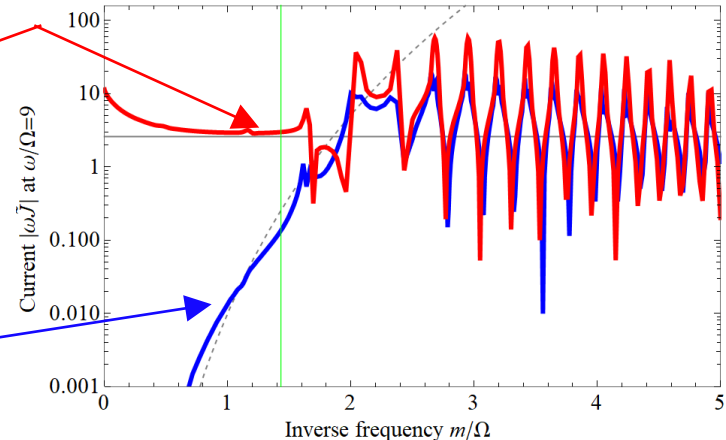


Dependence on strength E_0



Dependence on frequency Ω

analytical with WKB



numerical

• Lessons:

(1) WKB makes it easier to analyze the non-perturbative regime

(2) Saturation & oscillation of the harmonic intensity

⇒ consistent with recent semi-conductor exp. [Xia et al., (2020)]

(but only E_0 -dep. is measured and Ω -dep. is our prediction)

PART I: Basics of strong-field physics

PART II: Sauter-Schwinger effect and its connection to other areas

Summary

(QED = quantum electrodynamics)

Brief introduction to **Strong-field QED**

= an area to study **what happens by a super strong light**

(= electromagnetic field)

PART I: Basics of strong-field QED

- Why interesting \Rightarrow Unexplored non-perturbative regime of physics
Nontrivial phenomena that we've never seen e.g. Vacuum decay, Birefringence, New phase of matter, ...
- Why timely \Rightarrow Recent availability of strong fields e.g. High-power laser, Magnetar, heavy-ion collisions, ...

PART II: Sauter-Schwinger effect and its connection to other areas

- An interdisciplinary topic: Analogues in many other areas of physics
 \Rightarrow Exchange of ideas/techniques among physics is quite useful for better understandings
e.g. Landau-Zener (cond.-mat), Hawking radiation (gravity), reheating in the early Universe (cosmology), ...
- Mathematically, the problem is essentially solving a differential equation
 \Rightarrow Mathematical techniques to solve differential equations are quite useful
- As such an example: High-harmonic generation from the vacuum
 \Leftarrow Based on (exact) WKB in mathematics and an idea in cond.-mat
[HT, Fujimori, Misumi, Nitta, Sakai, (2021)] [HT, Hongo, Ikeda, (2021)]



Intuitive picture

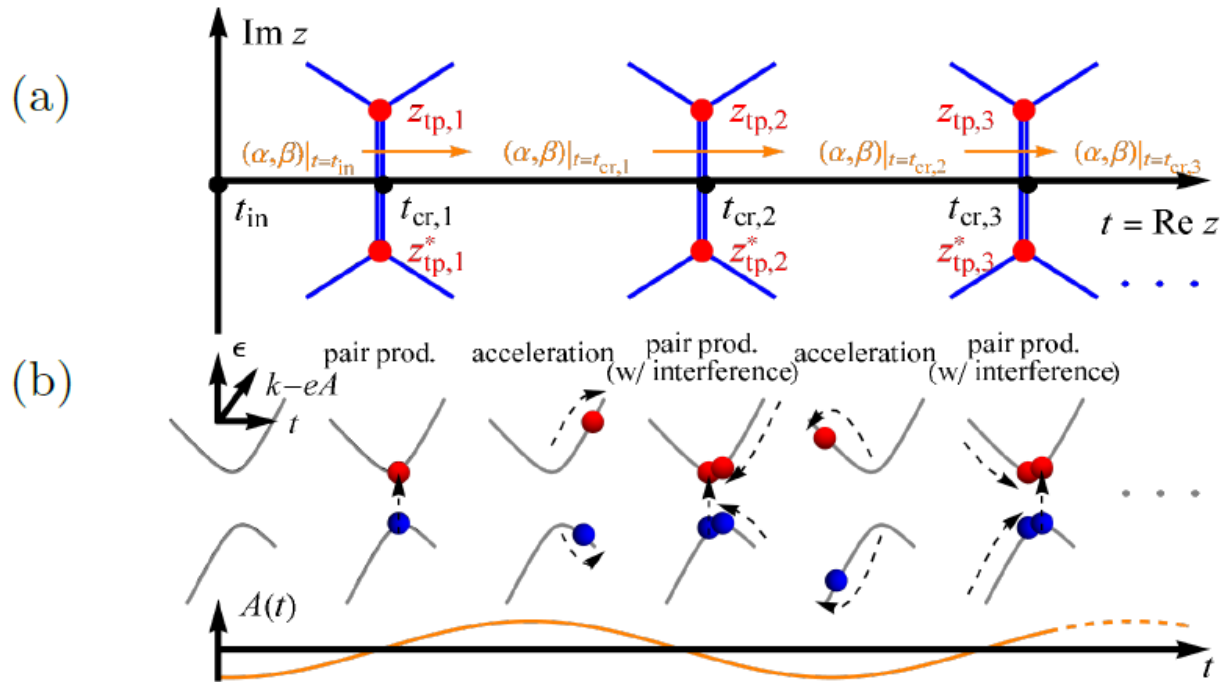


FIG. 1. (a) A typical Stokes graph, composed of Stokes lines (blue lines) and turning points (red points), and (b) the corresponding physical processes during the real-time evolution.