

Opportunities of strong-field physics in intermediate-energy heavy-ion collisions

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Plan

**Intermediate-energy heavy-ion collisions $\sqrt{s_{NN}} = O(2 - 10 \text{ GeV})$
is interesting not only to QCD/hadron but also to strong-field QED**

1. Introduction to strong-field QED

2. Strong EM field in high-energy heavy-ion collisions

Strong but too short-lived \Rightarrow affects “non-perturbativity” of strong-field processes

3. Strong EM field in intermediate-energy heavy-ion collisions

Estimate EM field profile with a hadron transport model (JAM)

\Rightarrow “strong” $O(50 \text{ MeV})$ and long-lived $O(10 \text{ fm}/c)$

\Rightarrow a nice setup to study strong-field QED; non-negligible to QCD/hadron processes as well

4. Summary

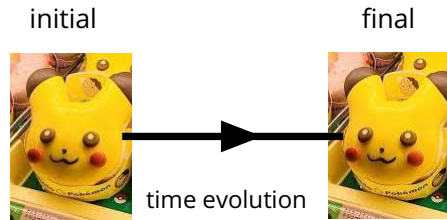
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2. Strong EM field in high-energy HIC

3. Strong EM field in intermediate-energy HIC

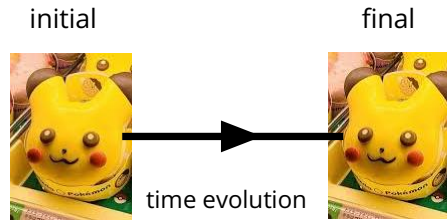
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Strong-field QED



Vacuum
(= No EM field)

Strong-field QED

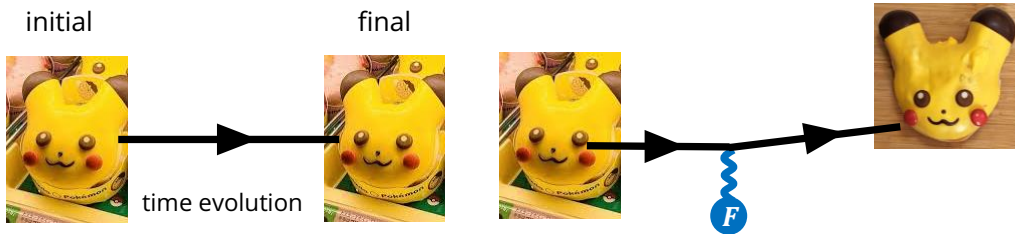


Vacuum
(= No EM field)

Weak EM field
($eF/m^2 \lesssim 1$)

Strong EM field
($eF/m^2 \gtrsim 1$)

Strong-field QED



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Strong EM field
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Almost the same

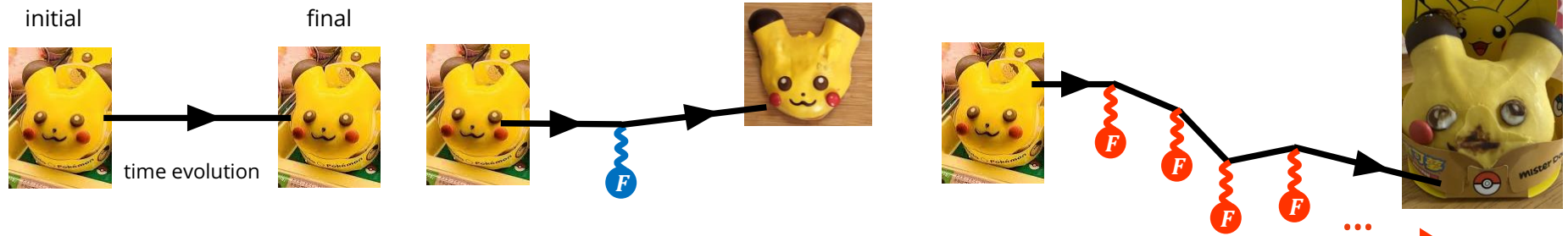
⇒ Perturbative

⇒ Understood

ex) Electron anomalous magnetic moment $a := \frac{g-2}{2}$

$$\begin{aligned} a(\text{theor.}) &= 1159652182.03 \dots \times 10^{-12} \\ a(\text{exp.}) &= 1159652180.73 \dots \times 10^{-12} \end{aligned} \quad [\text{Aoyama, Kinoshita, Nio (2017)}]$$

Strong-field QED



Vacuum
(= No EM field)

Weak EM field
($eF/m^2 \lesssim 1$)

Strong EM field
($eF/m^2 \gtrsim 1$)

Almost the same
⇒ Perturbative
⇒ Understood

Completely different
⇒ Non-perturbative
⇒ Not understood

ex) Electron anomalous magnetic moment $a := \frac{g-2}{2}$

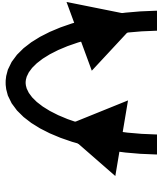
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Examples of strong-field phenomena

Novel QED processes ($eF/m_e^2 \gtrsim 1$)

Review: [Fedotov, Ilderton, Karbstein, King, Seipt, [HI](#), Torgrimsson (2022)]

ex) Schwinger effect



Photon splitting



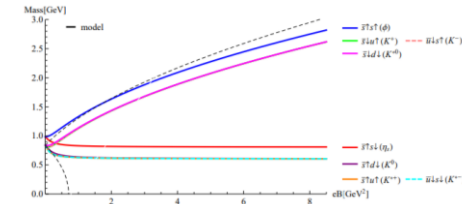
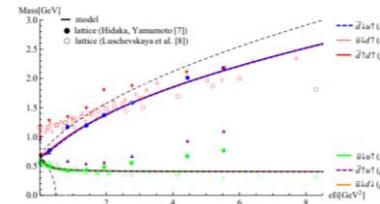
Vacuum birefringence
(= Polarization dep. of reflective index)



Impacts on QCD/hadron physics ($eF/\Lambda_{\text{QCD}}^2 \gtrsim 1$)

ex. 1) Hadron properties:

e.g., mass, charge dist., decay mode, ...



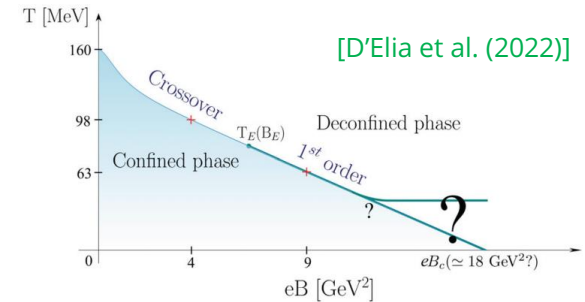
[[HI \(2015\)](#)]

ex. 2) QCD phase diagram

e.g., (inverse) magnetic catalysis, new phase, ...

ex. 3) Others: Anomalous transport,

(for color EM field) Glasma, string breaking, ...



[[D'Elia et al. \(2022\)](#)]

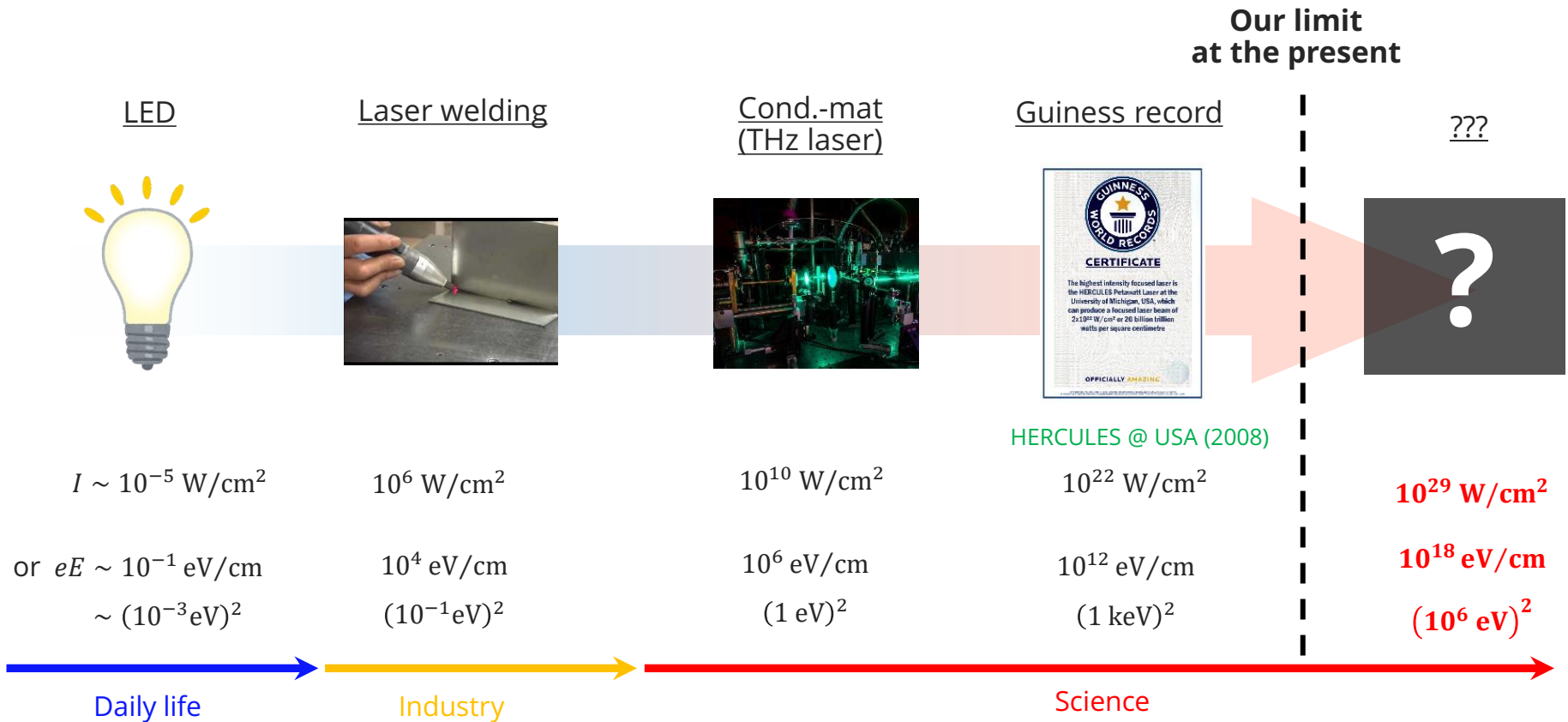
Many theoretical predictions, but NEVER observed in experiments

Need of EXTREMELY strong EM field

Order of the magnitude:

QED: $eE, eB > m_e^2 = (0.511 \text{ MeV})^2 \approx O(10^{29} \text{ W/cm}^2)$

QCD: $> \Lambda_{\text{QCD}}^2 = (200 \text{ MeV})^2 \approx O(10^{39} \text{ W/cm}^2)$



Impossible within the current tech. \Rightarrow New idea needed \Rightarrow HIC ?

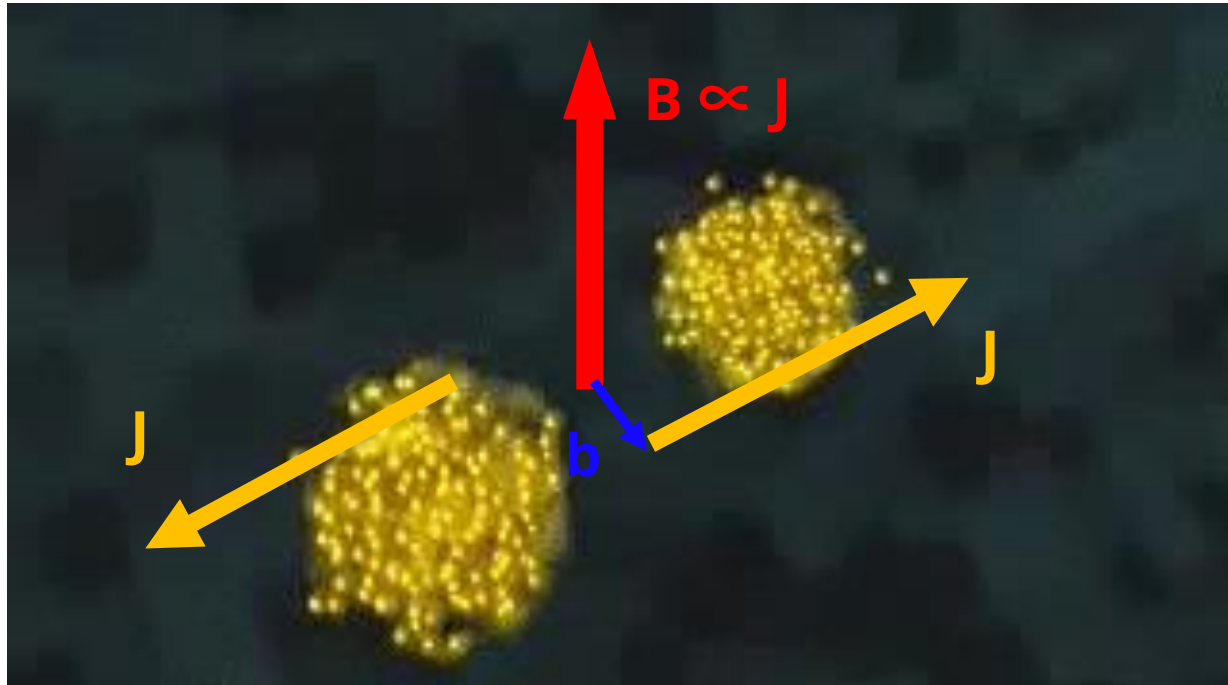
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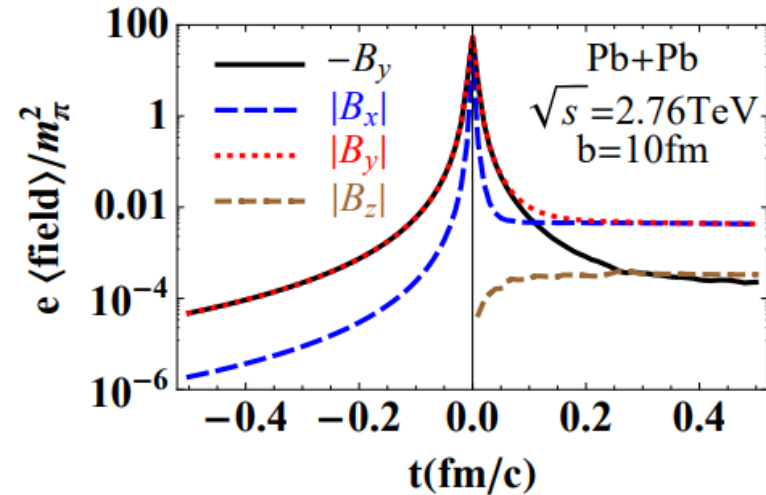
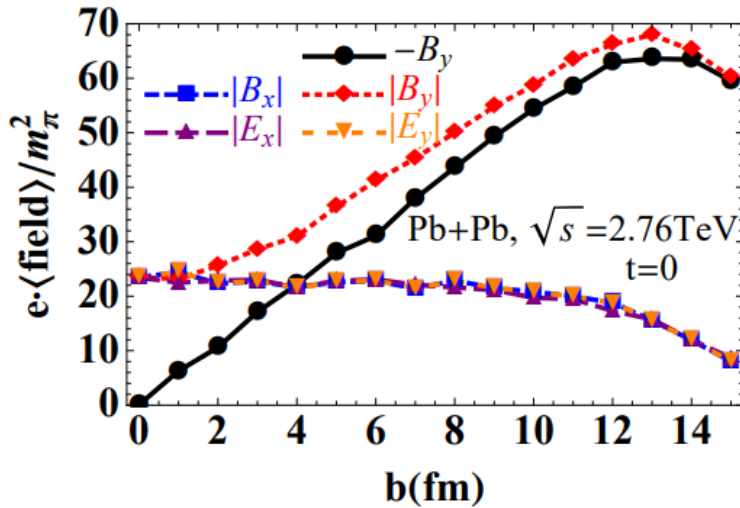
4. Summary

Strong EM field at high-energy HIC



- ✓ Strong magnetic field is created

Strong EM field at high-energy HIC



[Deng, Huang (2012)]

See also [Bzdak, Skokov (2012)] [Hattori, Huang (2016)]

✓ Strong magnetic field is created

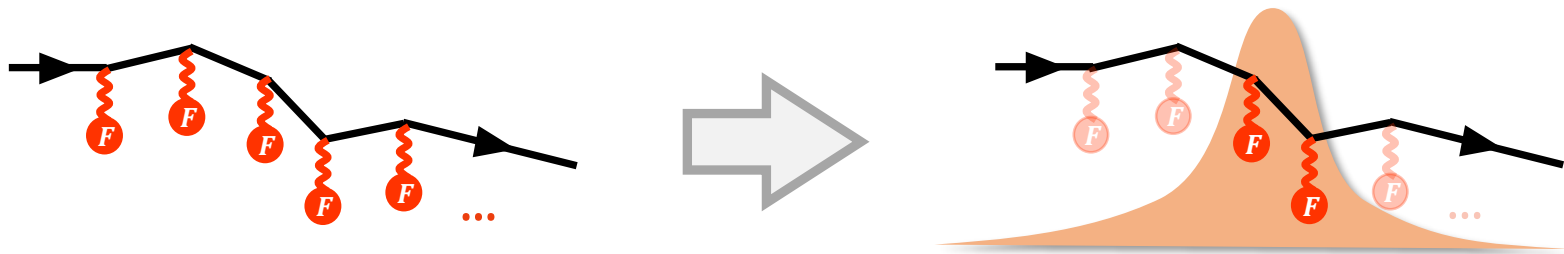
Pro: Super strong $eB \gg \Lambda_{\text{QCD}}^2$

Cons: Extremely short-lived ($\tau \ll 1$ fm/c)

⇒ Affects “non-perturbativity” of strong-field physics

Shorter lifetime \Rightarrow less non-perturbative

Intuition: No time for multiple interactions



Shorter lifetime \Rightarrow less non-perturbative

“Phase diagram” of strong-field physics

- As example: Vacuum particle prod. by E field w/ finite lifetime
- Three dimensional parameters in the system: eE, τ, m
 \Rightarrow Two dim.-less parameters determine the physics

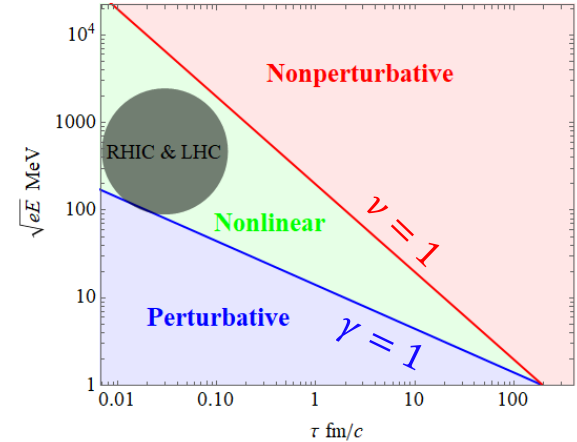
$$\gamma = \frac{m}{eE \tau} = \frac{\text{(Typical energy)}}{\text{(Work by field)}} \Rightarrow \text{Characterize the magnitude of work}$$

$$\nu = \frac{eE \tau}{1/\tau} = \frac{\text{(Work by field)}}{\text{(Photon energy)}} \Rightarrow \text{Characterize the number of photons}$$

- $\gamma \ll 1, \nu \gg 1 \Rightarrow$ Non-perturbative v.s. $\gamma \gg 1, \nu \ll 1 \Rightarrow$ perturbative

- High-energy HIC: $eF \sim (1 \text{ GeV})^2, \tau \sim 0.1 \text{ fm}/c \Rightarrow \gamma \sim \begin{cases} 10^{-3} (m = \Lambda_{\text{QCD}}) \\ 10^{-5} (m = m_e) \end{cases}, \nu \sim 0.1$

[HT, Fujii, Itakura (2014)]
 [HT, Fujimori, Misumi, Nitta, Sakai (2021)]



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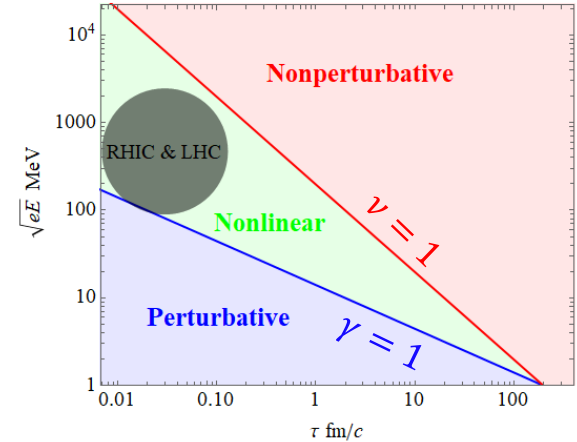
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High-energy heavy-ion collision is short-lived

\Rightarrow NOT useful for strong-field phys. in non-perturbative regime

Actually, only NLO processes such as Breit-Wheeler have been observed in exp.; no signals of higher-order effects

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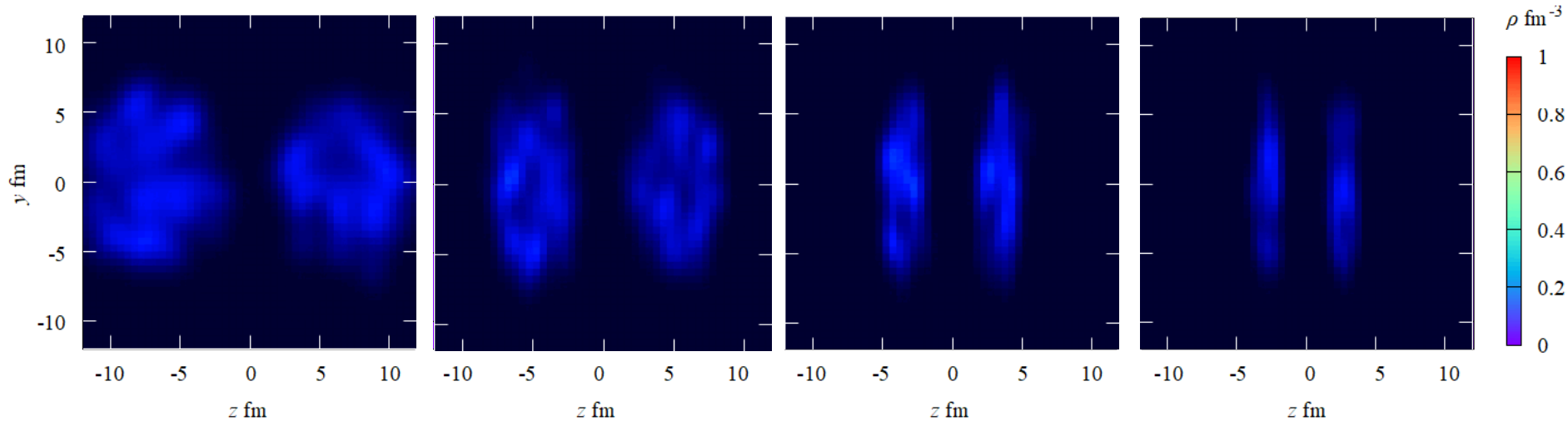
Intermediate-energy HIC

Heavy-ion collisions at $\sqrt{s_{NN}} = O(2 - 10)$

[AGS, SPS, RHIC BES, FAIR, NICA, HIAF, J-PARC-HI, ...]

$t = 0.0 \text{ fm}/c$ at $x = 0.0 \text{ fm}$

[HT, Jinno, Kitazawa, Nara (2024)]



$\sqrt{s_{NN}} = 2.4 \text{ GeV}$

3.9 GeV

7.2 GeV

11.5 GeV

Stopping
(Landau)



Penetration
(Bjorken)

- Idea: baryon stopping at lower energies \Rightarrow dense matter w/ long lifetime
- The lifetime can reach $O(10 \text{ fm}/c)$

Dense \Rightarrow Strong Coulomb field

Strong Coulomb electric field should be produced due to large $Z_{\text{tot}} = 2Z \sim 200 > \alpha^{-1}$

- Very rough order estimate

$$\text{Strength: } eE \sim \frac{Z\alpha}{r^2} \sim \Lambda_{\text{QCD}}^2 \sim (100 \text{ MeV})^2$$

$$\text{Lifetime: } \tau \sim 10 \text{ fm}/c$$

$$\Rightarrow \gamma = \frac{m}{eE\tau} \lesssim \begin{cases} 10^{-1} (m = \Lambda_{\text{QCD}}) \\ 10^{-4} (m = m_e) \end{cases} \sim 0.1, \nu = eE\tau^2 \gtrsim 10$$

\Rightarrow Non-perturbative $\begin{cases} \gamma \ll 1 \\ \nu \gg 1 \end{cases}$ both for QED & QCD !

- If this is true, it's super interesting, since this is the very first physical system where we can study strong-field physics in the non.-pert regime
- But, of course, it's too rough, so let's do a realistic estimate


Approach: Hadronic transport model JAM

[Nara, Otsuka, Ohnishi, Nitta, Chiba (2000)]

- A successful model to simulate the realtime dynamics of heavy-ion collisions, reproducing various data (v1, yields, ...)
- Basic idea: superposition of collisions of individual hadrons (incl. inelastic channels such as resonance, string breaking, mini-jet)
- JAM returns the distribution of charged hadrons at each spacetime pt.

⇒ EM fields can be obtained as

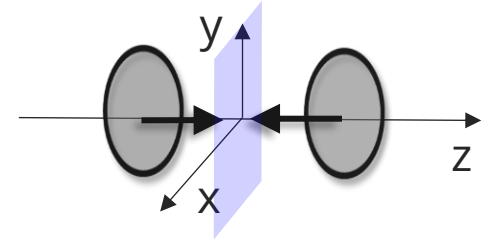
$$A^\mu(x^0, \mathbf{x}) = \frac{1}{4\pi} \int_{-\infty}^{+\infty} d^3\mathbf{x}' \frac{J^\mu(x^0 - |\mathbf{x} - \mathbf{x}'|, \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|}$$

From JAM


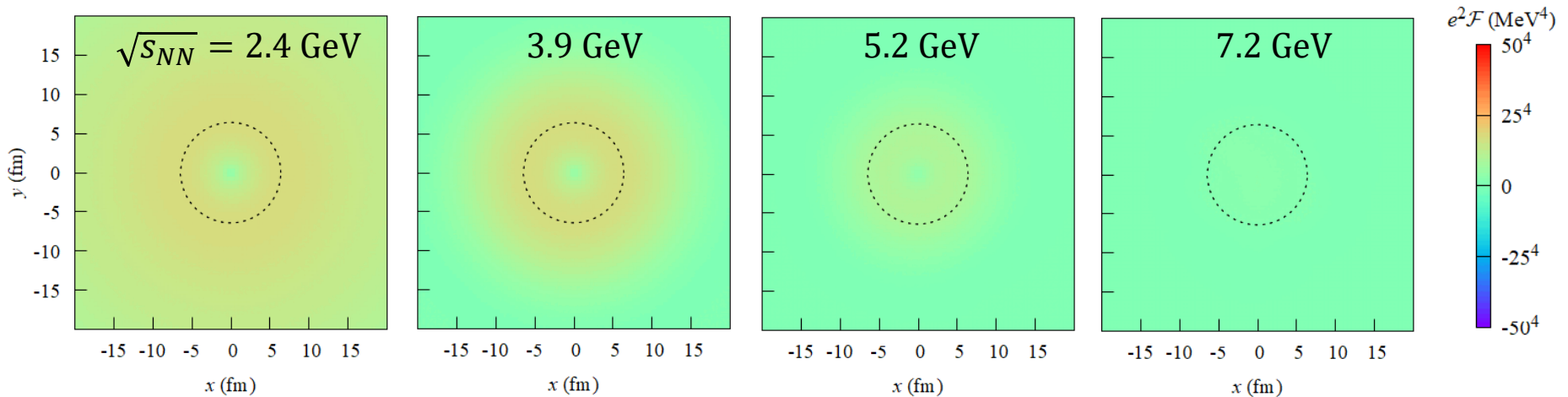
- NB: Just one of the models, not a first-principle calculation (e.g., no quark/gluon DoGs, no hydro, no phase transition, ...)
 - ⇒ should be regarded as a “baseline”, before incl. non-trivial physics
 - ⇒ worth to compare w/ other models: UrQMD, HIJING, SMASH, ...

Result (1/5): Spacetime profile at central coll.

✓ Event-averaged $F := E^2 - B^2$ ($F > 0$: Electric, $F < 0$: Magnetic)



$t = 0.0 \text{ fm}/c$ at $z = 0.0 \text{ fm}$



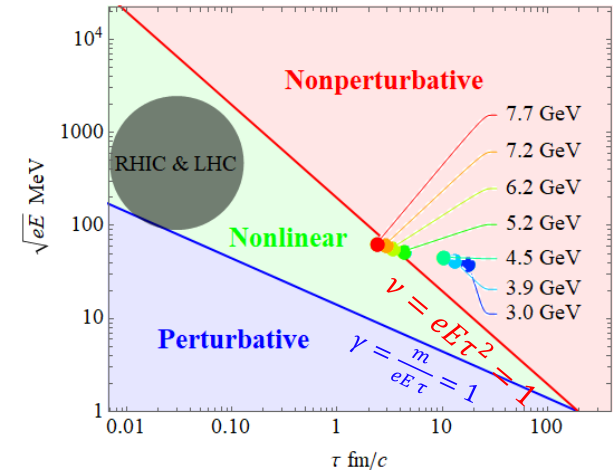
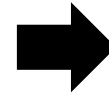
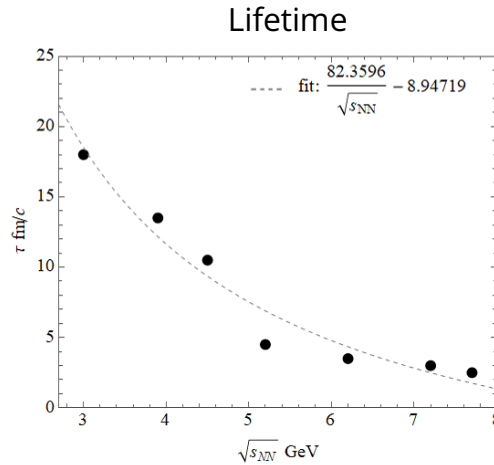
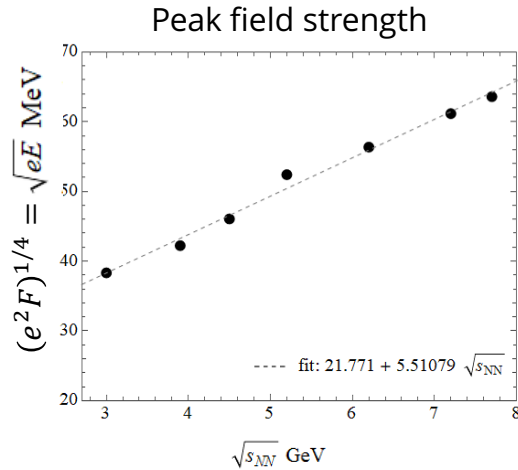
[HT, Nishimura, Ohnishi (2024)]

- F is positive \Rightarrow E field dominates over B field
- Donuts shaped \Leftarrow Gauss law $E \propto \int d^3x \rho$
- “strong” $O(50 \text{ MeV})$ and long-lived $O(10 \text{ fm}/c)$

↖ much stronger than m_e , current quark mass; non-negligible to $m_\pi, \Lambda_{\text{QCD}}$

Result (2/5): Strength & lifetime at central coll.

✓ Peak field strength and lifetime (FWHM)

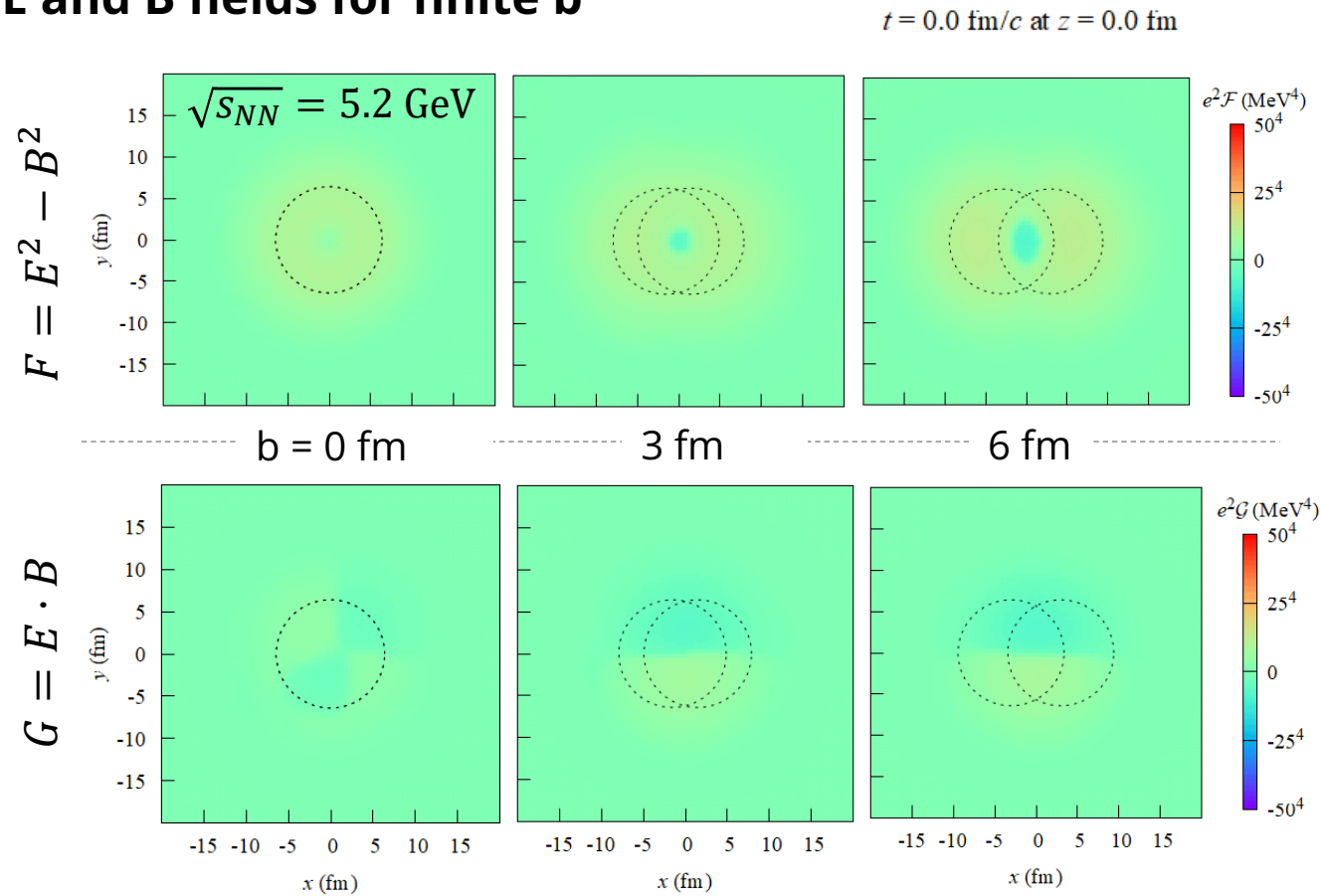
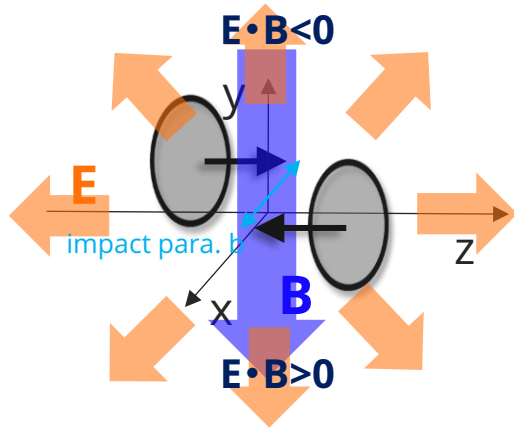


[HT, Nishimura, Ohnishi (2024)]

- Two basic physics: Lorentz contraction, Baryon stopping
- Intermediate energies can explore the non-perturbative regime
 - ⇐ Long lifetime compensates the weakness of the field

Result (3/5): Spacetime profile at **non-central coll.**

✓ Interplay between E and B fields for finite b



- B field appears but E field is always larger in space

[HT, in progress]

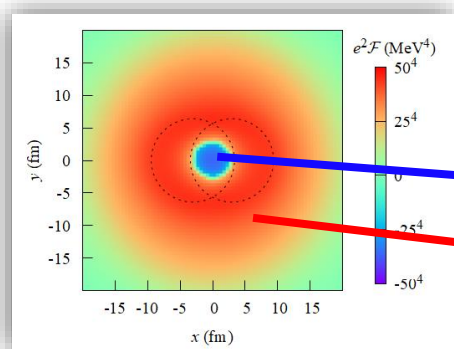
⇒ E field would be more important than B field in intermediate energies

- Parallel EM configuration s.t. $G = E \cdot B \neq 0$

⇒ can be a source of chiral physics $\partial_\mu J_5^\mu \propto E \cdot B$

Result (4/5): Strength & lifetime of **F** at **non-central coll.**

[HT, in progress]



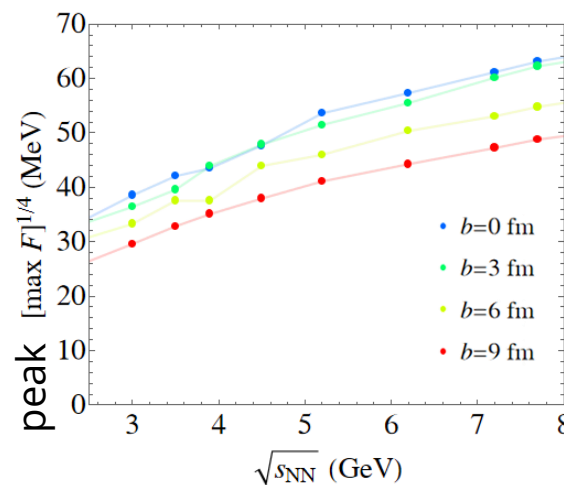
- E field is always strong compared to B field

⇒ E field is important even at peripheral events for intermediate energies

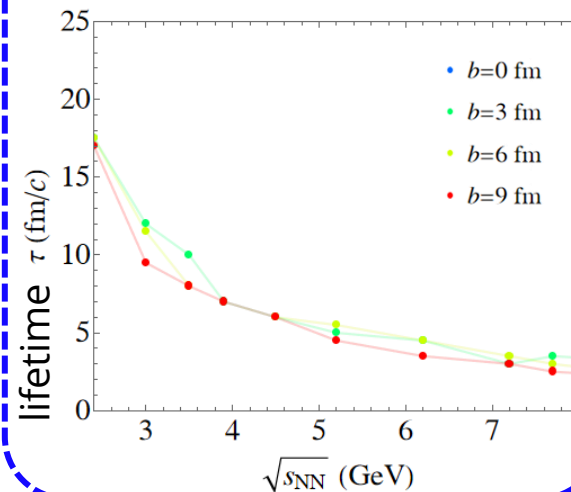
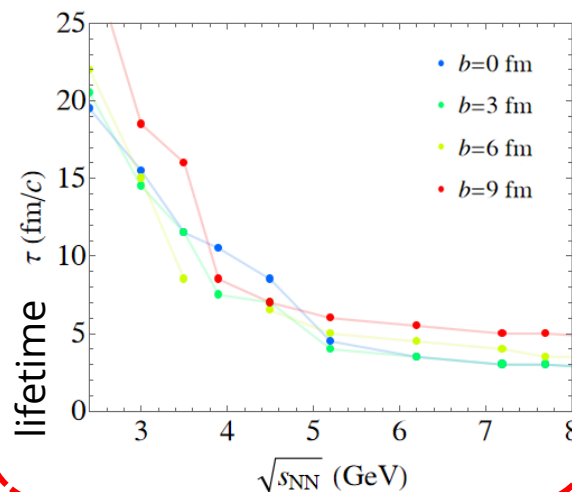
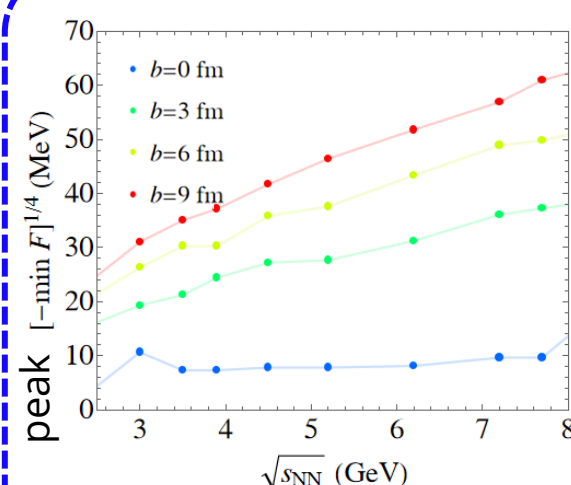
- Lifetimes are roughly the same and less sensitive to impact parameter b .

⇒ E & B fields can be non-pert. even at peripheral events

max $F \sim E$ field

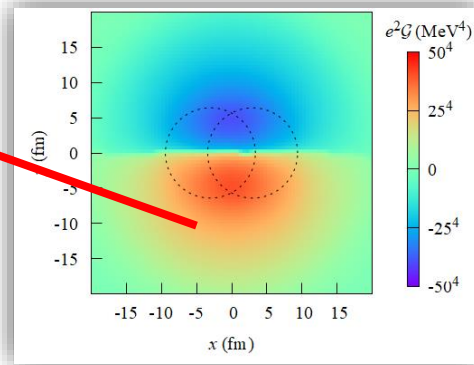
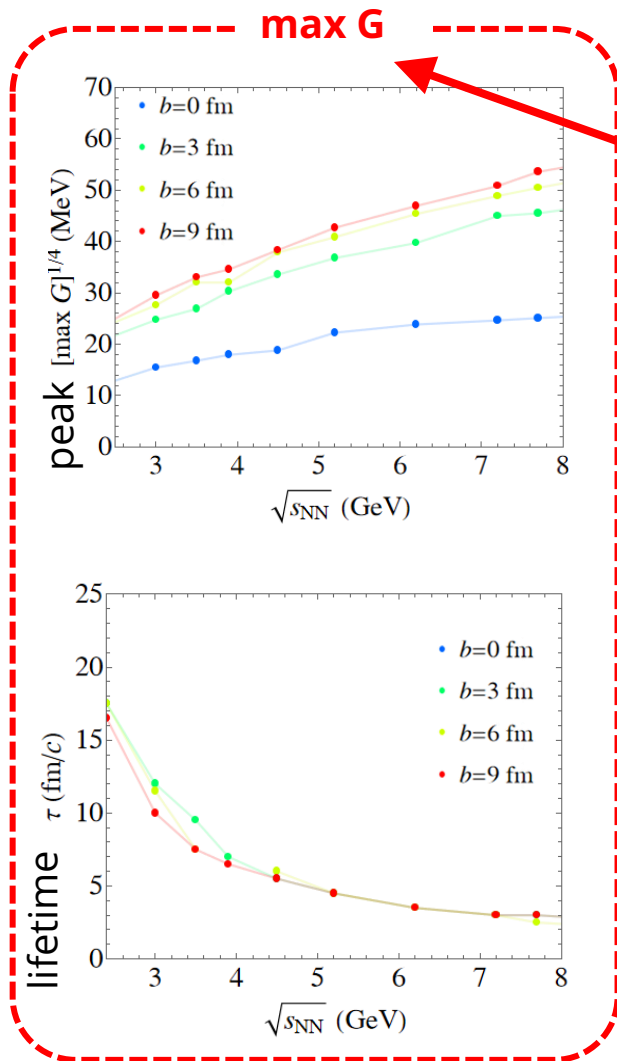


min $F \sim B$ field



Result (5/5): Strength & lifetime of **G** at **non-central coll.**

[HT, in progress]



- “Strong” as $G \sim F = O(50 \text{ MeV})$
 - As long-lived as F as $\tau = O(10 \text{ fm}/c)$
- ⇒ Not a simple E or B field configuration, but need to think both F and G effects in intermediate energy heavy-ion collisions

(similar plot for max G \doteq — min G)

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- Strong but too short-lived \Rightarrow affects “non-perturbativity” of strong-field processes

✓ Strong EM field in intermediate-energy heavy-ion collisions

- Estimate EM field profile with a hadron transport model (JAM)
- Coulomb electric field is produced, which is “strong” $\mathcal{O}(50 \text{ MeV})$ and long-lived $\mathcal{O}(10 \text{ fm}/c)$
 \Rightarrow a novel opportunity to study strong-field QED; non-negligible to QCD/hadron as well
- E field is more important than B field
- “Chiral” $E \cdot B \neq 0$ configuration in peripheral collisions

